

Suppressing and restoring the Dicke superradiance transition by dephasing

Jonathan Keeling & Peter Kirton



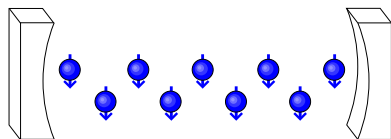
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Madrid, January 2017

Dicke model and Dicke-Hepp-Lieb transition

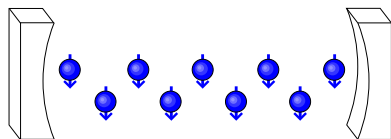


$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \frac{\omega_0}{2} \sigma_{\alpha}^z + g(\psi + \psi^\dagger)(\sigma_{\alpha}^+ + \sigma_{\alpha}^-)$$

- Coherent state: $|\Psi\rangle \rightarrow e^{\lambda \psi^\dagger + \eta S^+} |\Omega\rangle$
- Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

Dicke model and Dicke-Hepp-Lieb transition

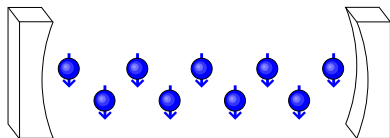


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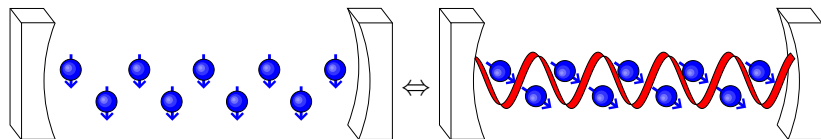
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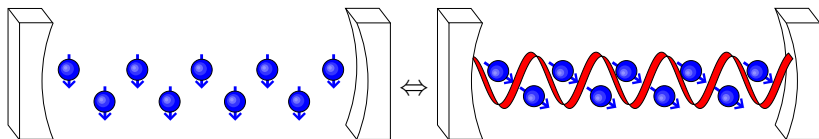
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Non-zero cavity field if: $4Ng^2 > \omega\omega_0$

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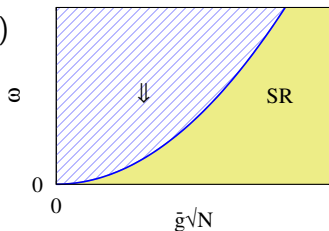
Dicke model and Dicke-Hepp-Lieb transition



$$\begin{aligned}
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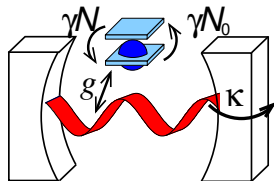
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“Textbook” Laser: Maxwell Bloch equations

$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \frac{\omega_0}{2} \sigma_{\alpha}^z + g \left(\psi \sigma_{\alpha}^+ + \psi^\dagger \sigma_{\alpha}^- \right)$$



- $|\psi\rangle^2 > 0$ if $2N_0g^2 > \gamma n$
- Requires inversion

“Textbook” Laser: Maxwell Bloch equations

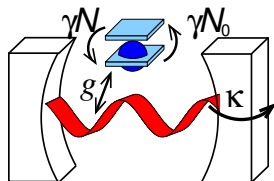
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Maxwell-Bloch eqns: $P = -i\langle\sigma^-\rangle$, $N = 2\langle\sigma^z\rangle$

$$\partial_t\psi = -i\omega\psi - \frac{\kappa}{2}\psi + \sum_{\alpha} gP_{\alpha}$$

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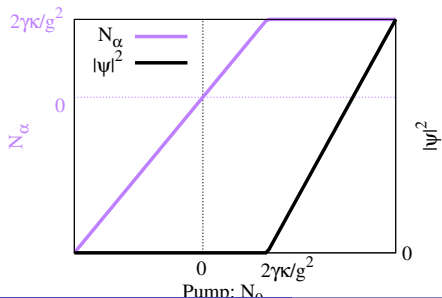
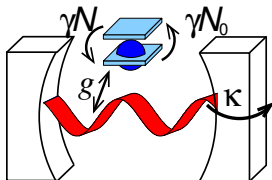
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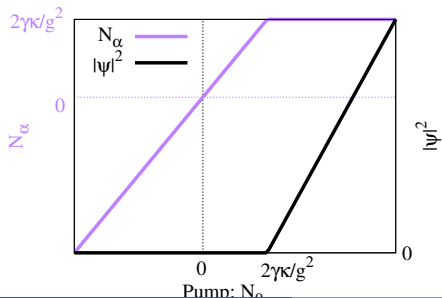
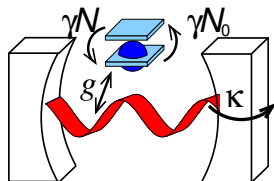
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“Textbook” Laser: Semiclassical equations

- Semiclassical laser theory $n = \langle \psi^\dagger \psi \rangle$

$$\partial_t n = \gamma N_0 \frac{2g^2(n+1)}{\gamma\gamma_t + 4g^2(n+1)} - \kappa n$$

- MF Transition at $N_0 2g^2 / \gamma_t = \kappa$
- No symmetry breaking
- Spontaneous emission: finite “size” corrections

$$n = \frac{1}{2\beta} \left[\frac{N_0}{N_c} - 1 \pm \sqrt{\left(\frac{N_0}{N_c} - 1 \right)^2 + 4\beta \frac{N_0}{N_c}} \right]$$

$$\beta = 4g^2 / \gamma\gamma_t$$

[Haken, RMP, 1975]

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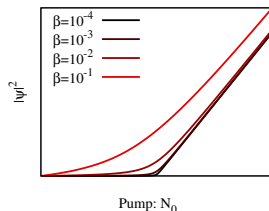
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1 Introduction: Open Dicke model reminder

- Dicke superradiance vs lasing
- Driven Dicke model

2 Behaviour with dephasing

- Mean field theory problem
- Polynomial algorithm for exact solution
- Cumulant expansion and $N \rightarrow \infty$

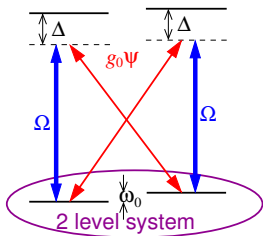
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Raman driven Dicke model



$$H = \omega_0 S^Z + g(\psi + \psi^\dagger)(S^- + S^+) + \omega \psi^\dagger \psi$$

- 2 Level system, $|\downarrow\rangle, |\uparrow\rangle$

- Coupling $g = \frac{g_0 \Omega}{2\Delta}$

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- Imbalanced case (internal states):

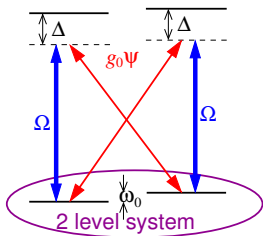
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[Dimer *et al.* PRA '07]

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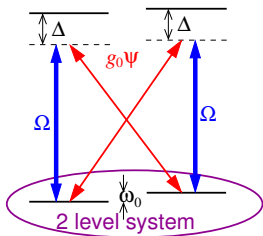
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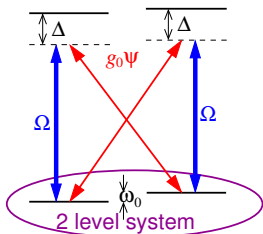
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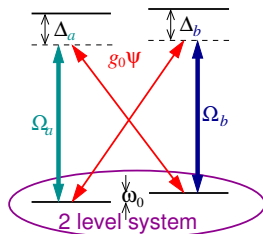
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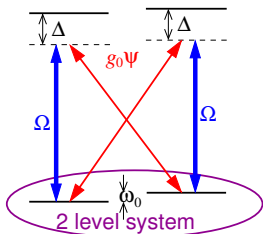
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[Dimer et al. PRA '07]

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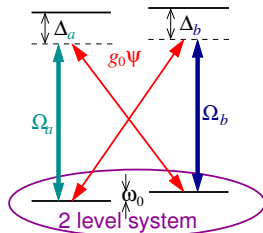
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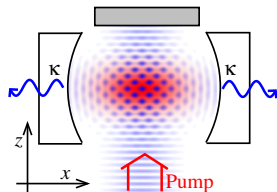
Open Dicke model theory

- Momentum degrees of freedom:

$$\psi = \psi_{\downarrow} + \psi_{\uparrow} \cos(kx) \cos(kz)$$

- Effective 2LS ($\psi_{\downarrow}, \psi_{\uparrow}$)

$$H_{\text{eff}} = \omega \psi^{\dagger} \psi + \sum_n \frac{\omega_0}{2} \sigma_n^z + g_{\text{eff}} \sigma_n^x (\psi + \psi^{\dagger}) + U \sigma_n^z \psi^{\dagger} \psi$$



- Extra “feedback” term U , cavity loss κ

• Single mode – mean-field EOM, $\alpha = \langle \psi \rangle$, $S^i = \sum_n \sigma_n^i / 2$.

$$S^{\pm} = -i(\omega_0 + U|\alpha|^2) S^{\pm} + 2g_{\text{eff}}(\alpha + \alpha^*) S^{\pm}$$

$$S^z = ig_{\text{eff}}(\alpha + \alpha^*)(S^{\pm} - S^{\mp})$$

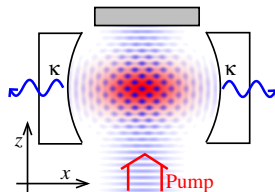
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$$\dot{\alpha} = -[\kappa + i(\omega + US^z)] \alpha - ig_{\text{eff}}(S^- + S^+)$$

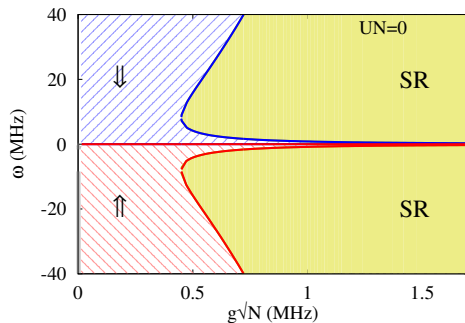
Classical dynamics

Changing U :

$$U = 0$$

$$U < 0$$

$$U > 0$$



[JK *et al.* PRL '10, Bhaseen *et al.* PRA '12]

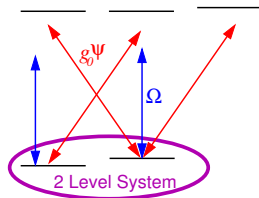
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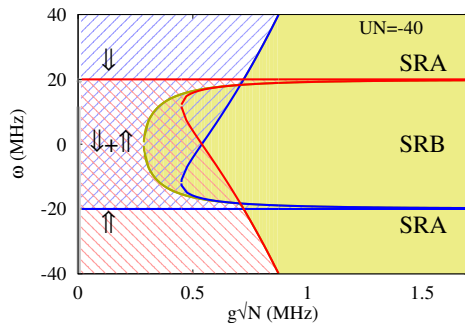
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$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$



[JK *et al.* PRL '10, Bhaseen *et al.* PRA '12]

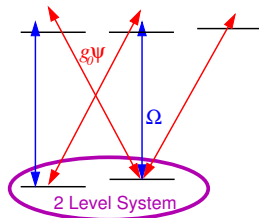
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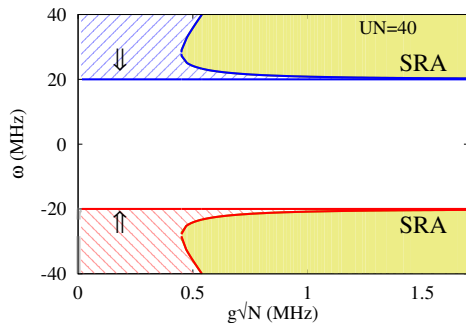
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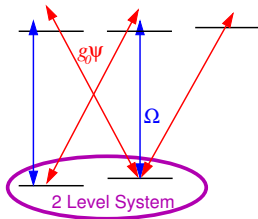
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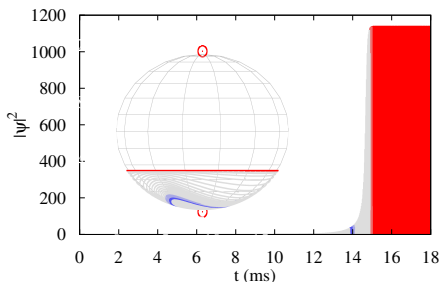
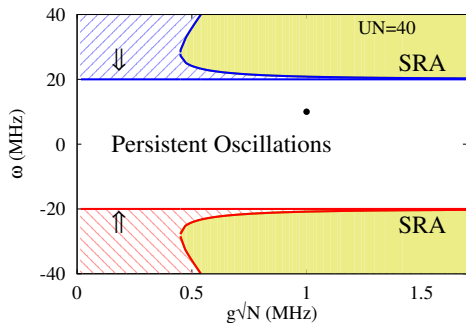
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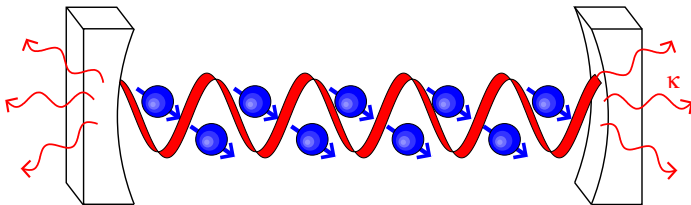
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Individual dephasing



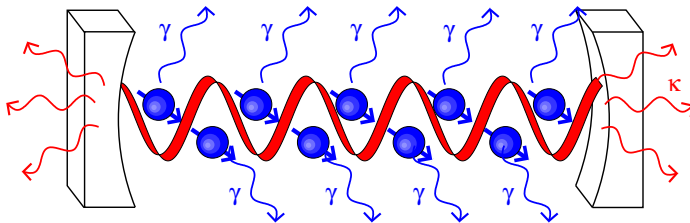
• Extra loss terms

$$\dot{\rho}_{ij} = -i[H, \rho] + \kappa \mathcal{D}[\rho] + \sum_a \Gamma_a \mathcal{D}[\rho_a] + \Gamma_b \mathcal{D}[\rho_b]$$

$$\mathcal{L}[\rho] = X\rho X^\dagger - (X^\dagger X\rho + \rho X^\dagger X)/2$$

• Γ_a, Γ_b break S conservation.

Individual dephasing



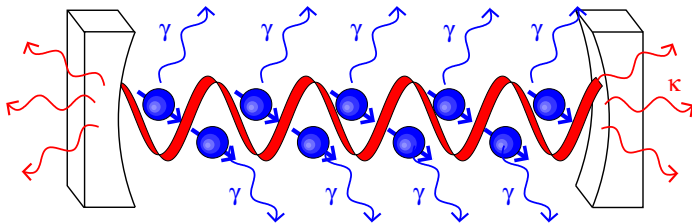
• Extra loss terms

$$a_{pp} = -i(H_p) + \kappa Q[V] + \sum_{\alpha} \Gamma_{\alpha} Q[\rho_{\alpha}] + \Gamma_{\alpha} Q[\rho_{\alpha}^{\dagger}]$$

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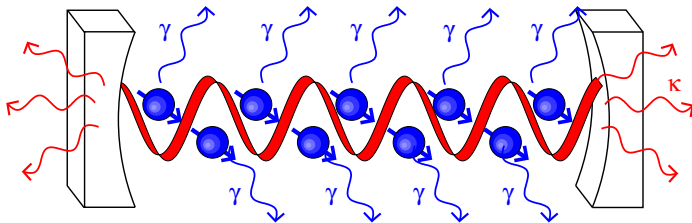
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Effect of incoherent losses

- Adding other loss terms

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$$\partial_t \langle \psi \rangle = -i\omega \langle \psi \rangle - igN \langle \sigma^z \rangle - \kappa \langle \psi \rangle / 2$$

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- Steady state and Linear stability

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- Steady state *and* Linear stability

Maxwell-Bloch: limiting cases

$$\begin{aligned}
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$$2g^2 N > \frac{-1}{\langle \sigma^z \rangle} \left(\frac{\omega^2 + (\kappa/2)^2}{\omega} \right) \left(\frac{\omega_0^2 + (\Gamma_\phi + \Gamma_\downarrow/4)^2}{\omega_0} \right)$$

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Exact solution

- No fixed $\mathbf{S} = \sum_i \sigma_i$: $\text{Dim}[\mathcal{H}] = 2^N$ vs N

• But: Permutation symmetry of ρ remains

• Evolve projected ρ : size $N^4 \times (\Omega_{\text{phot,max}})^2$.

[Kirton & JK, arXiv:1611.03342]

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→ $R(\dots s_n \dots s_m \dots) = R(\dots s_m \dots s_n \dots)$

→ Need only ordered list of $0 \leq \epsilon_i < 4$

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Exact solution

Wigner function $W(\psi = x + ip)$,

- Finite N : no symmetry breaking
 - Superradiance: bimodal state
- Γ_0 only unimodal
- Suggestive, inconclusive:

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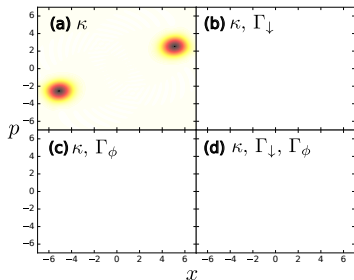
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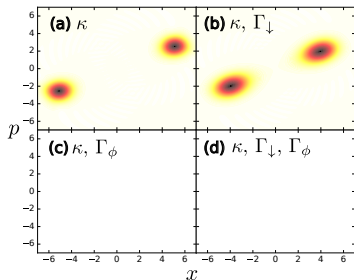


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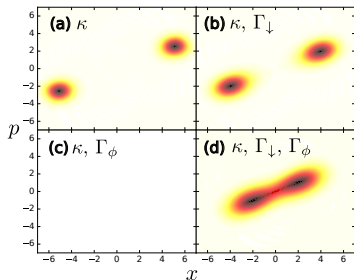


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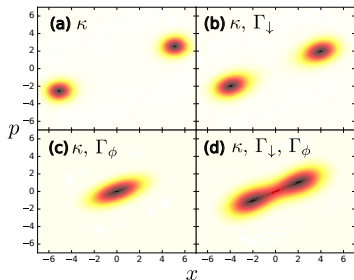
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Asymptotic behaviour

- Proof of transition: Finite size scaling

- ▶ Superradiant:

- $\langle \psi^\dagger \psi \rangle \propto N$

- ▶ Normal: $\langle \psi^\dagger \psi \rangle \propto \sqrt{N}$

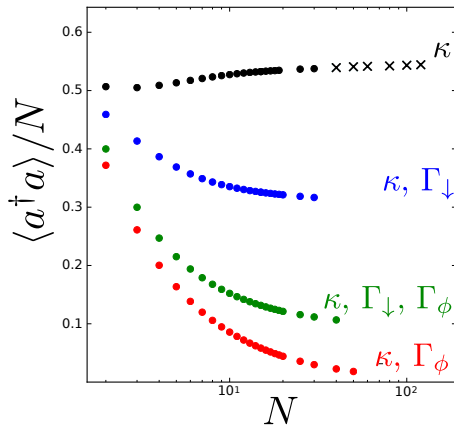
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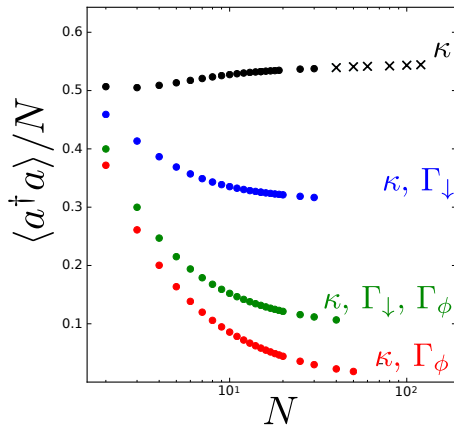


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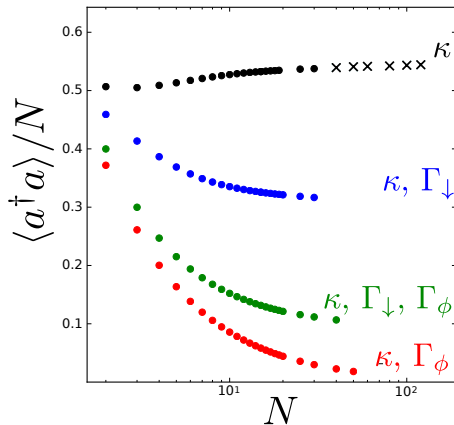
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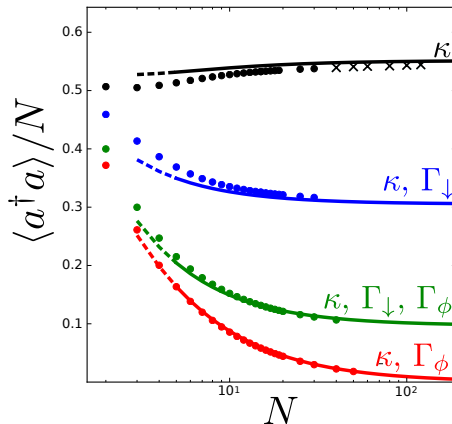
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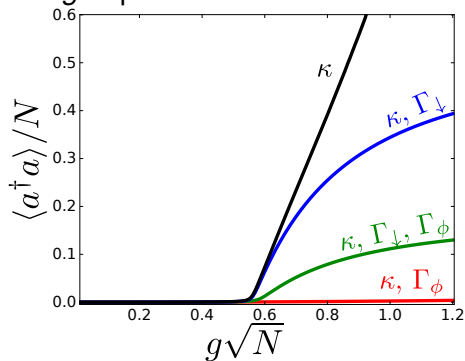
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Further applications

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$$H = \dots + g \psi^{\dagger} \sigma_{\alpha}^{-} + g' \psi^{\dagger} \sigma_{\alpha}^{+} + \text{H.c.} + \dots$$

• $g' > 0, \Gamma_{\uparrow} = 0$, driven Dicke

• $g' = 0, \Gamma_{\uparrow} > 0$, Laser

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- $g' > 0, \Gamma_{\uparrow} = 0$, driven Dicke

- $g' = 0, \Gamma_{\uparrow} > 0$, Laser

Further applications

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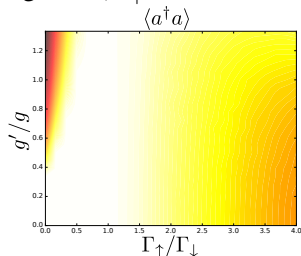
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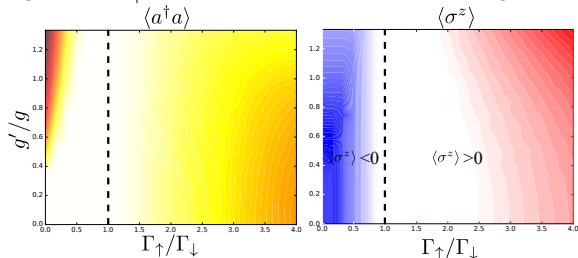


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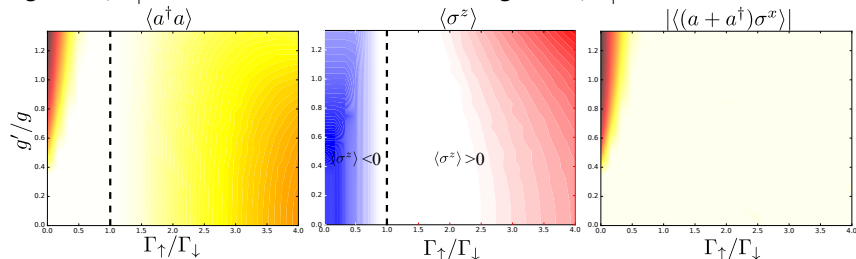
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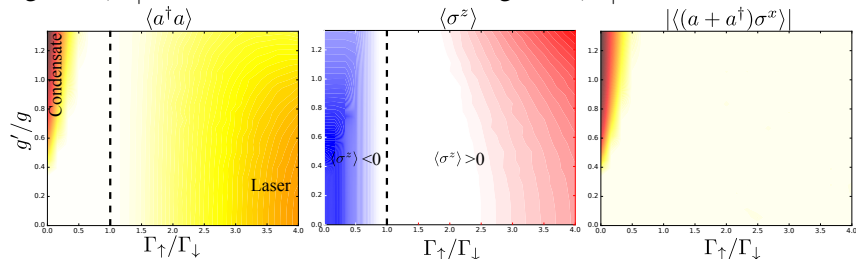
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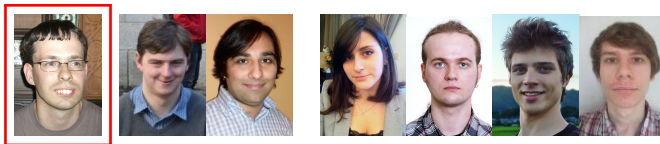
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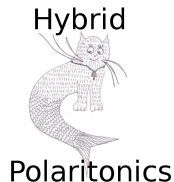


Acknowledgements

GROUP:

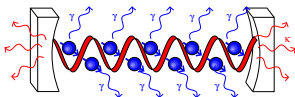


FUNDING:

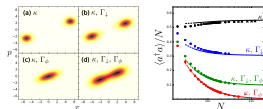


Summary

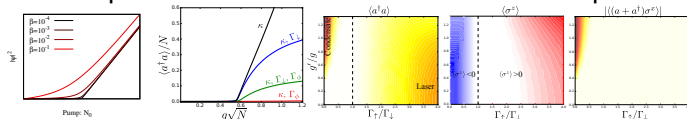
- Open Dicke model, $\kappa, \Gamma_\phi, \Gamma_\downarrow$



- Exact numerics: permutation symmetry of ρ



- Cumulant expansion — connection to laser rate equation



[Kirton & JK, arXiv:1611.03342]

Green's function as common language

- Green's function: Response to weak perturbation

$$\left[D^R(\nu) \right]^{-1} = \nu - \omega + i\frac{\kappa}{2} + \frac{g^2 N_0}{\nu - \omega_0 + i\Gamma}$$

- Normal modes
- Linear stability

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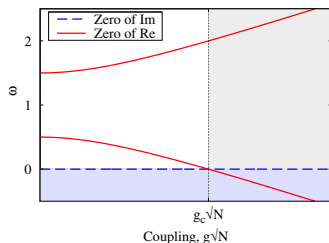
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Ground-state transition



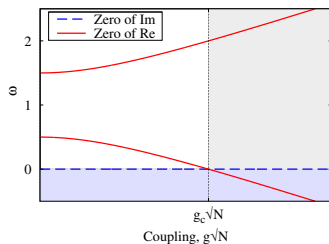
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Laser

