

Non-equilibrium phases of coupled matter-light systems

Jonathan Keeling



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St Andrews

600
YEARS



Southampton, May 2013

Coupling many atoms to light

Old question: *What happens to radiation when many atoms interact “collectively” with light.*

Superradiance — dynamical and steady state.

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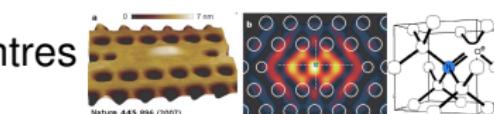
Superradiance — dynamical and steady state.

New relevance

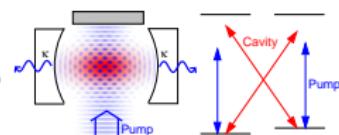
- Superconducting qubits



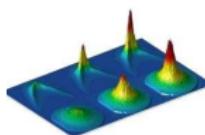
- Quantum dots & NV centres



- Ultra-cold atoms



- Rydberg atoms/polaritons



- Microcavity Polaritons

Dicke effect: Enhanced emission

PHYSICAL REVIEW

VOLUME 93, NUMBER 1

JANUARY 1, 1954

Coherence in Spontaneous Radiation Processes

R. H. DICKE

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey



$$H_{\text{int}} = \sum_{k,i} g_k (\psi_k^\dagger S_i^- e^{-i\mathbf{k}\cdot\mathbf{r}_i} + \text{H.c.})$$

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If $|\mathbf{r}_i - \mathbf{r}_j| \ll \lambda$, use $\sum_i \mathbf{S}_i \rightarrow \mathbf{S}$
Collective decay:

$$\frac{d\rho}{dt} = -\frac{\Gamma}{2} [S^+ S^- \rho - S^- \rho S^+ + \rho S^+ S^-]$$

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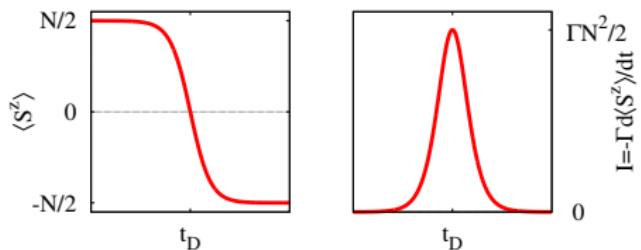


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$$I \propto -\Gamma \frac{d\langle S^z \rangle}{dt} = \frac{\Gamma N^2}{4} \operatorname{sech}^2 \left[\frac{\Gamma N}{2} t \right]$$



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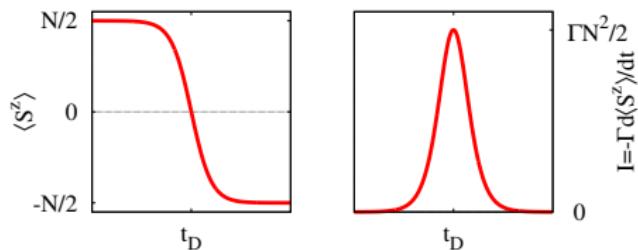


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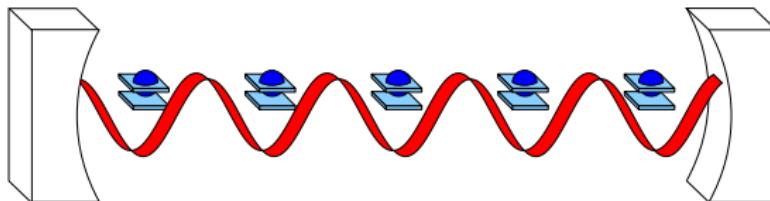
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Problem: dipole interactions dephase. [Friedberg et al, Phys. Lett. 1972]

Collective radiation **with** a cavity: Dynamics

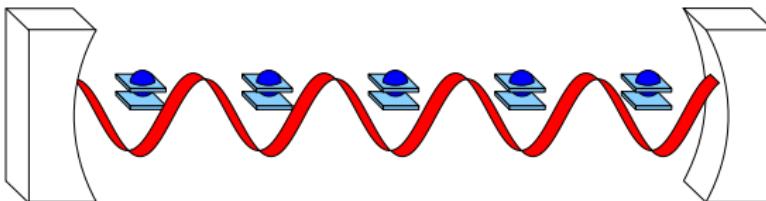


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Single cavity mode: oscillations

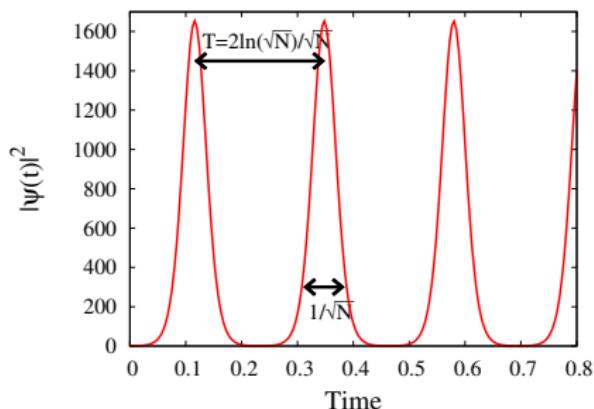
[Bonifacio and Preparata PRA '70]

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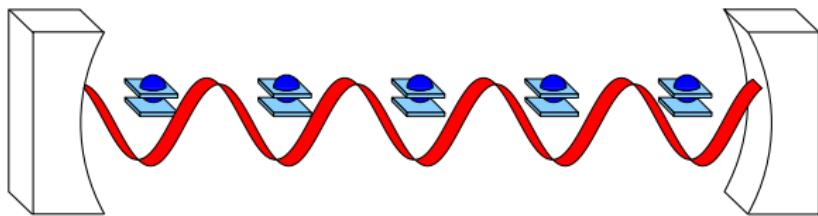
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Dicke model: Equilibrium superradiance transition



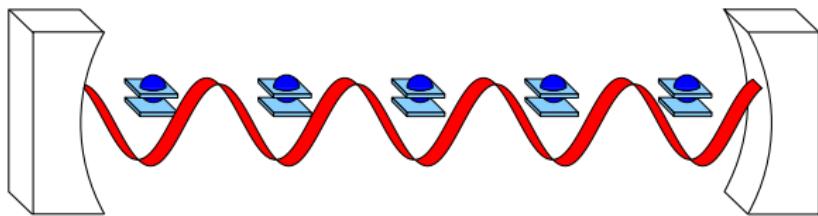
$$H = \omega\psi^\dagger\psi + \omega_0 S^z + g(\psi^\dagger S^- + \psi S^+).$$

• Ground state (Poissonian)

• Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

Dicke model: Equilibrium superradiance transition



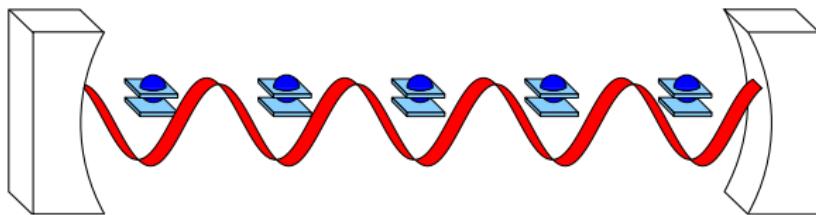
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- Coherent state: $|\Psi\rangle \rightarrow e^{\lambda\psi^\dagger + \eta S^+} |\Omega\rangle$

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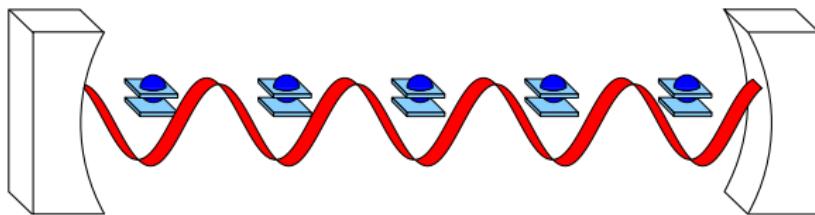
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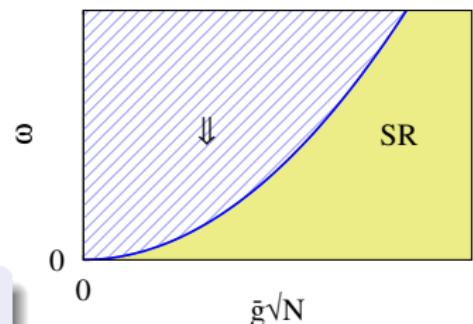
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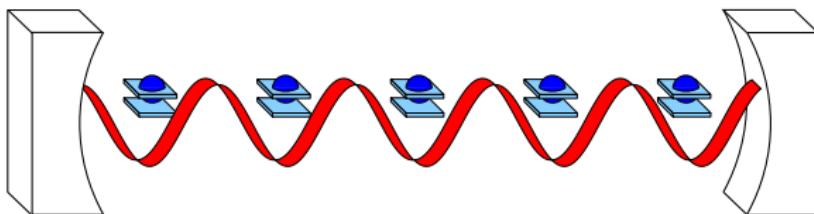
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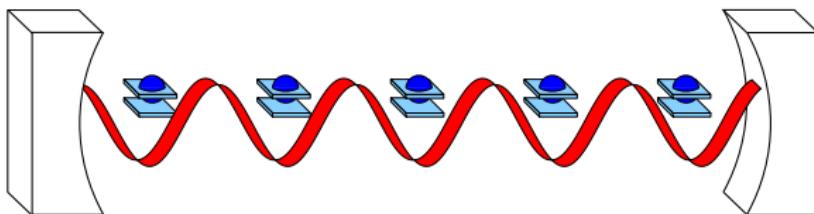
No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\omega_0$

[Rzazewski *et al* PRL '75]

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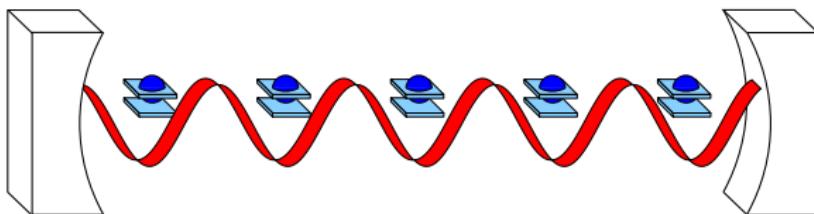
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$$-\sum_i \frac{e}{m} A \cdot p_i \Leftrightarrow g(\psi^\dagger S^- + \psi S^+), \quad \sum_i \frac{A^2}{2m} \Leftrightarrow N\zeta(\psi + \psi^\dagger)^2$$

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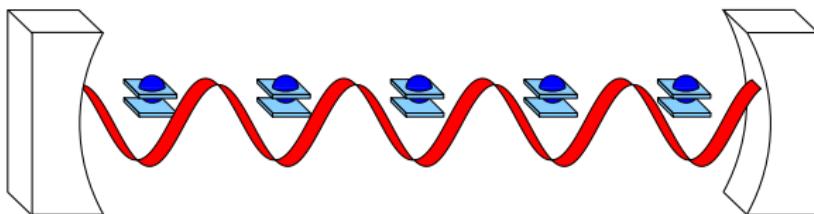
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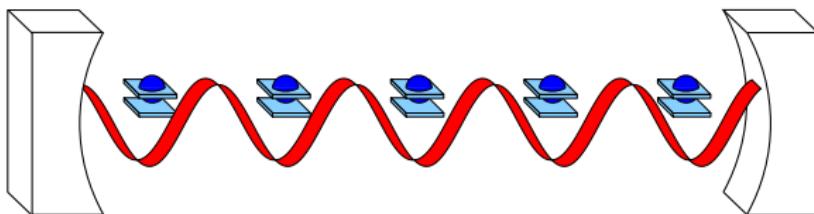
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But Thomas-Reiche-Kuhn sum rule states: $g^2/\omega_0 < 2\zeta$. **No transition**
[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation:
Ferroelectric transition in D+ γ gauge.
[Kapchinskii, Volosov & Demokritov PRA 2012]
- Circuit QED [Nataf and Ciuti, Nat. Comm. '10; Viehmann et al. PRL '11]
- Grand canonical ensemble:
 - If $H \rightarrow H - \mu(5^+ + 5^-)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
 - Incoherent pumping — polariton condensation.
- Dissociate g, ω_0 ,
e.g. Raman scheme: $\omega_R \ll \omega$.
(Dimer et al. PRA '07; Baumann et al. Nature '10; Also, Black et al. PRL '03)

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Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.

[JK JPCM '07, Vukics & Domokos PRA 2012]

- Change field parameter and coupling: $\omega_0 \gg \mu$, [Mehmann et al. PRL '11]
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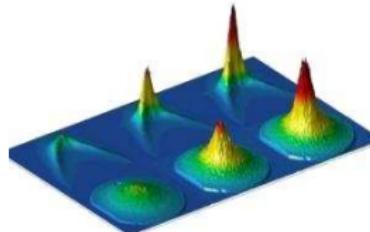
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→ Dicke phase transition

→ e.g., Raman scattering $\omega < \omega_0$

(Dimer *et al.* PRA '07; Baumann *et al.* Nature '10. Also, Black *et al.* PRL '09)

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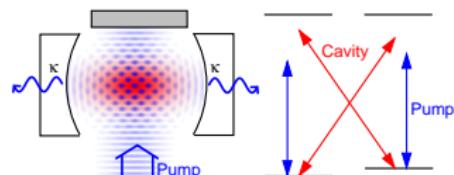
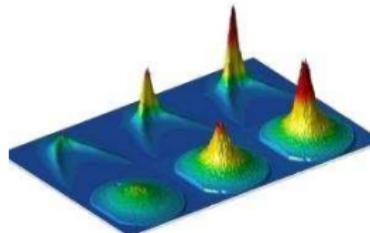
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Outline

1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- Timescales and consequences for experiment
- Persistent oscillating phases

3 Jaynes Cummings Hubbard model

- JCHM vv Dicke
- Coherently driven array
- Disorder

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GROUP:



COLLABORATORS: Simons, Bhaseen, Schmidt, Blatter, Türeci, Krüger

EXPERIMENT: Houck, Wallraff, Fink, Mylnek

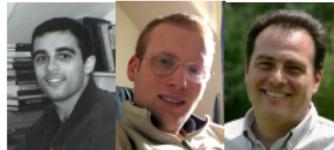
FUNDING:



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Research Council



Dynamics of generalized Dicke model



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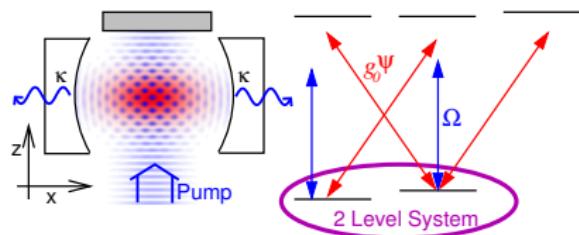
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Reminder of cold-atom extended Dicke model



2 Level system, $|\Downarrow\rangle, |\Uparrow\rangle$:

$$\Downarrow: \Psi(x, z) = 1$$

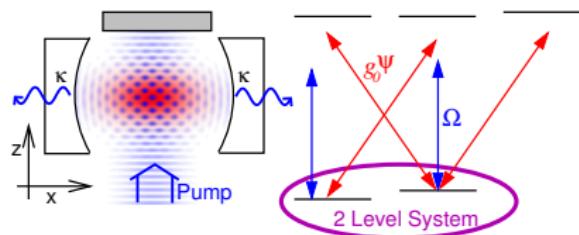
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$$H = \omega \psi^\dagger \psi + \omega_0 S^z + g(\psi + \psi^\dagger)(S^- + S^+) - i S_\perp \omega \psi$$

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[Baumann et al Nature '10]

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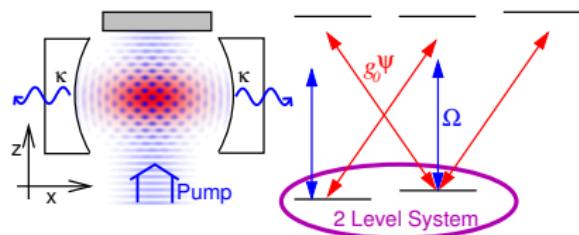
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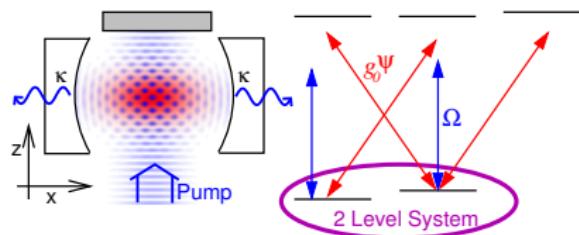
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Feedback: $\textcolor{red}{U} \propto \frac{g_0^2}{\omega_c - \omega_a}$

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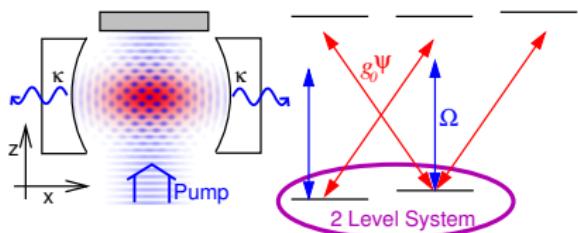
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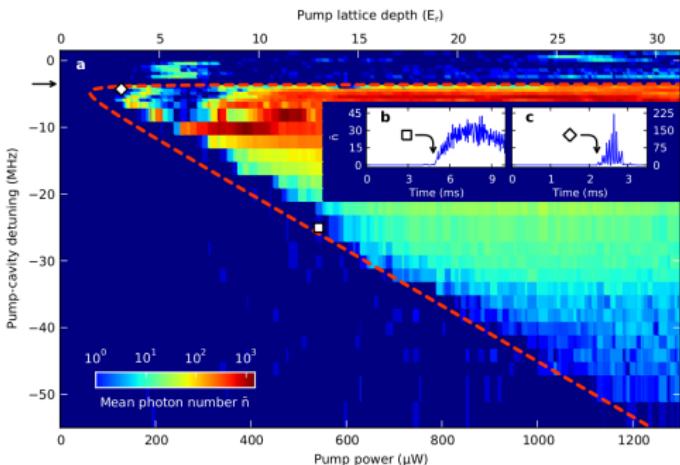


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Classical dynamics of the extended Dicke model

Open dynamical system:

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- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
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Classical EOM
 $(|\mathbf{S}| = N/2 \gg 1)$

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- Neglect quantum fluctuations → classical mechanics for large N , small κ
- initial conditions
- linearisation about fixed point
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→ Equilibrium classical fixed points

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Open dynamical system:

$$H = \omega\psi^\dagger\psi + \omega_0 S^z + g(\psi + \psi^\dagger)(S^- + S^+) + US_z\psi^\dagger\psi.$$

$$\partial_t\rho = -i[H, \rho] - \kappa(\psi^\dagger\psi\rho - 2\psi\rho\psi^\dagger + \rho\psi^\dagger\psi)$$

Classical EOM
 $(|\mathbf{S}| = N/2 \gg 1)$

$$\begin{aligned}\dot{S}^- &= -i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z \\ \dot{S}^z &= ig(\psi + \psi^*)(S^- - S^+) \\ \dot{\psi} &= -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)\end{aligned}$$

- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - ▶ Recover Retarded Green's function (spectrum)
 - ▶ Cannot recover occupations

Fixed points (steady states)

$$0 = i(\omega_0 + \textcolor{red}{U}|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

$\psi = 0, S = (0, 0, \pm N/2)$

$$0 = ig(\psi + \psi^*)(S^- - S^+)$$

$\Rightarrow \text{if } g > g_c, \psi \neq 0 \text{ too}$

$$0 = -[\kappa + i(\omega + \textcolor{red}{U}S^z)]\psi - ig(S^- + S^+)$$

$$\begin{cases} S^z = -g[S^+] = 0 \\ \psi = g[S^+] = 0 \end{cases}$$

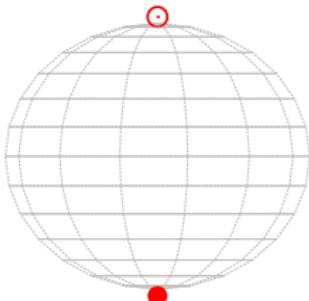
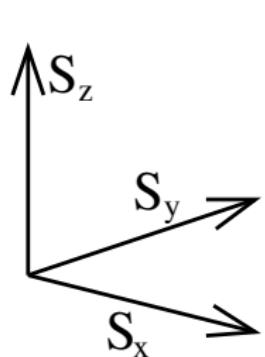
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- $\psi = 0, \mathbf{S} = (0, 0, \pm N/2)$ always a solution.



Small g: \uparrow, \downarrow only.
($\omega = 30\text{MHz}$, $UN = -40\text{MHz}$)

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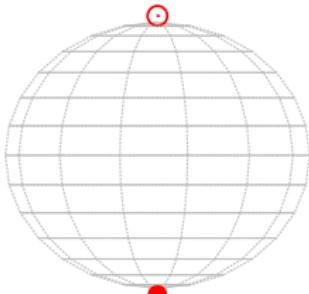
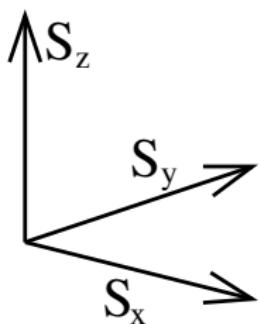
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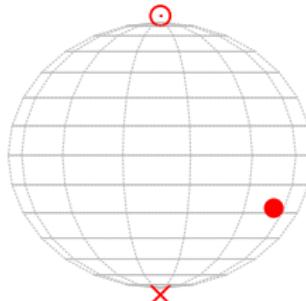
- $\psi = 0, \mathbf{S} = (0, 0, \pm N/2)$ always a solution.

- If $g > g_c$, $\psi \neq 0$ too

- A $S^y = -\Im[S^-] = 0$
- B $\psi' = \Re[\psi] = 0$



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($\omega = 30\text{MHz}$, $UN = -40\text{MHz}$)



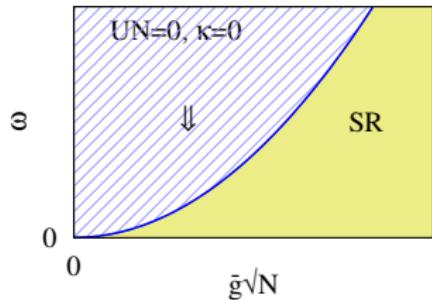
Larger g: SR too.

Steady state phase diagram

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

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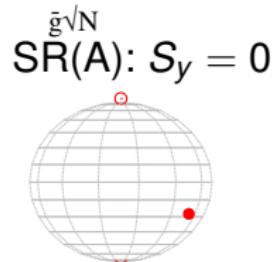
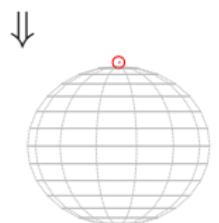
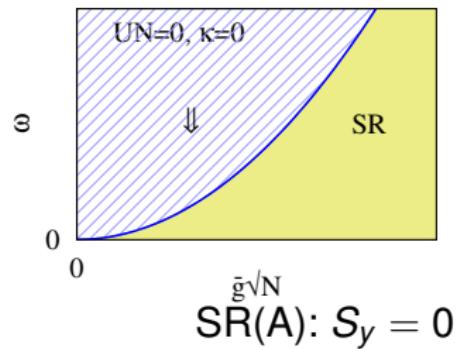
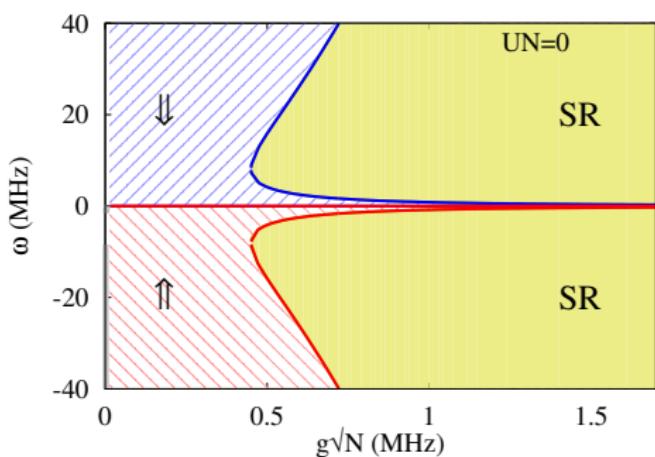
See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

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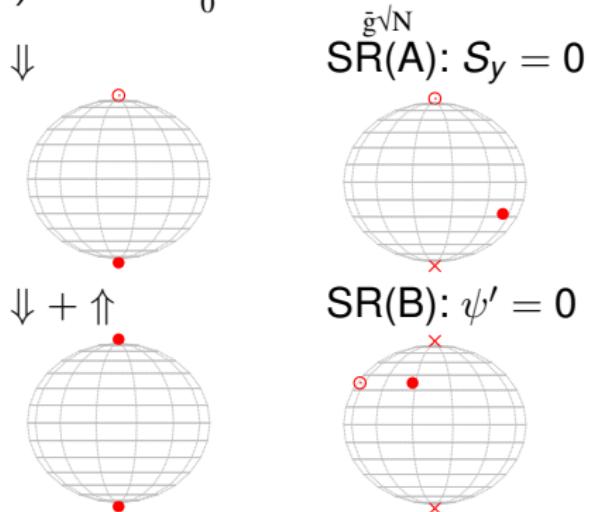
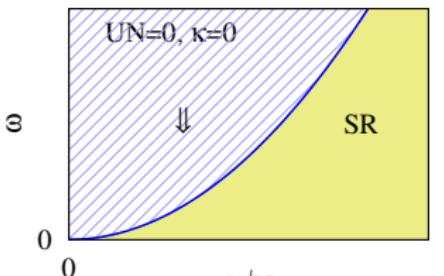
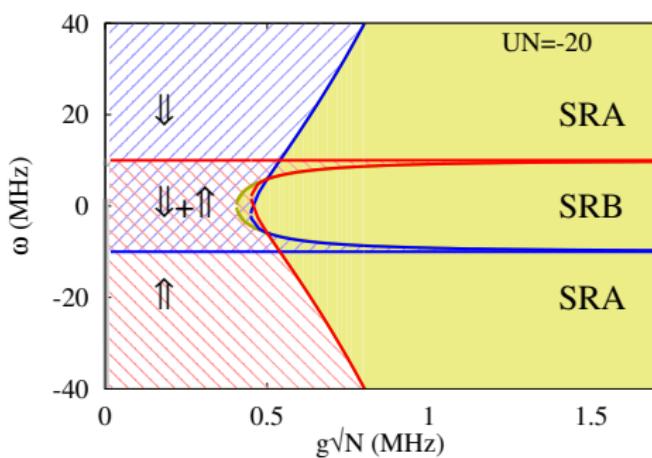
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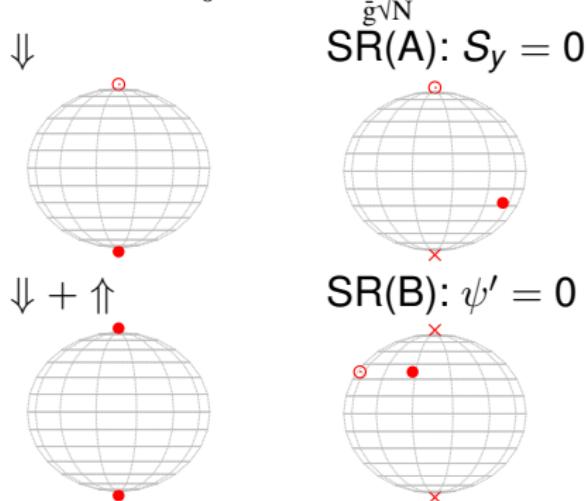
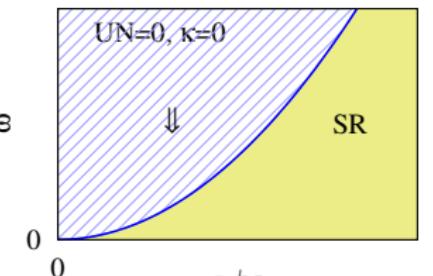
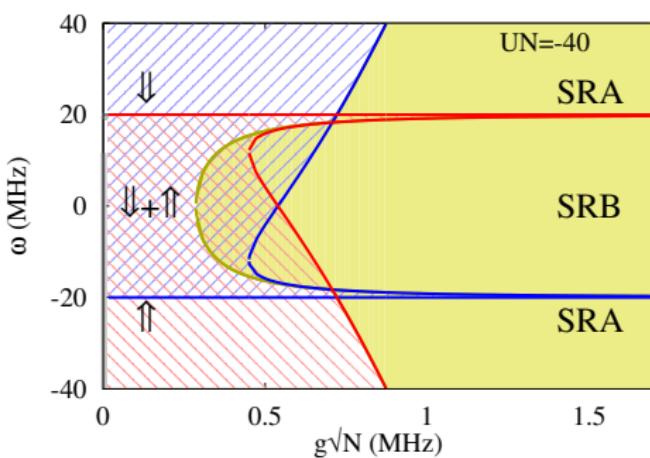
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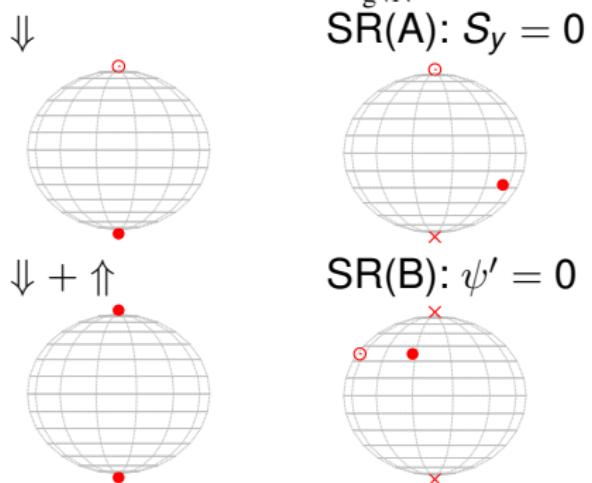
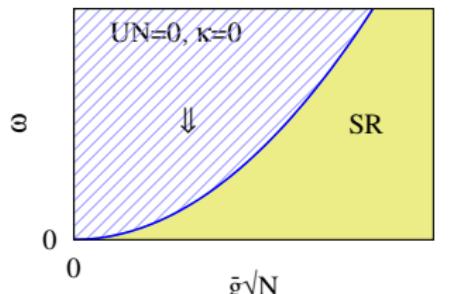
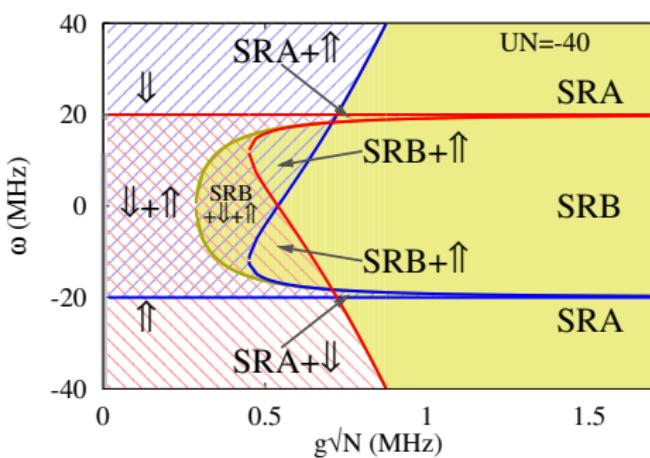
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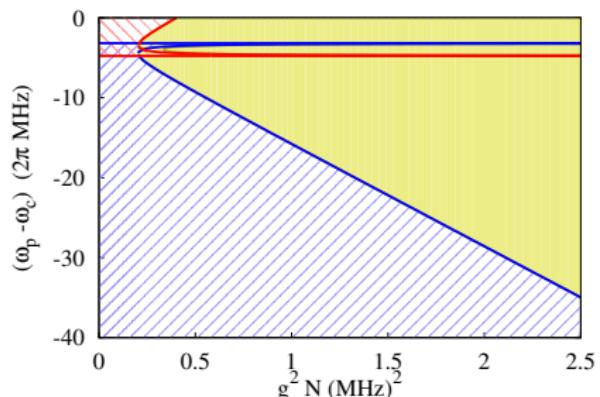
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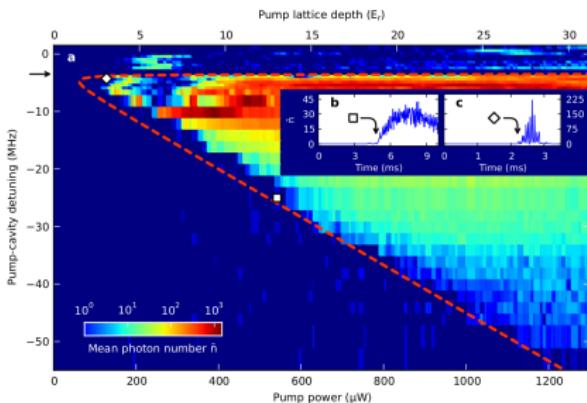
Comparison to experiment



$$UN = -10 \text{ MHz}$$

Adapted from: [Bhaseen *et al.* PRA '12]

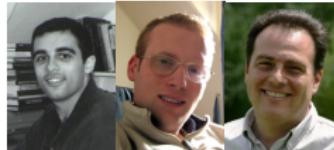
$$\omega = \omega_c - \omega_p + \frac{5}{2} UN,$$



[Baumann *et al* Nature '10]

$$UN = -\frac{g_0^2}{4(\omega_a - \omega_c)}$$

Dynamics of generalized Dicke model



1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

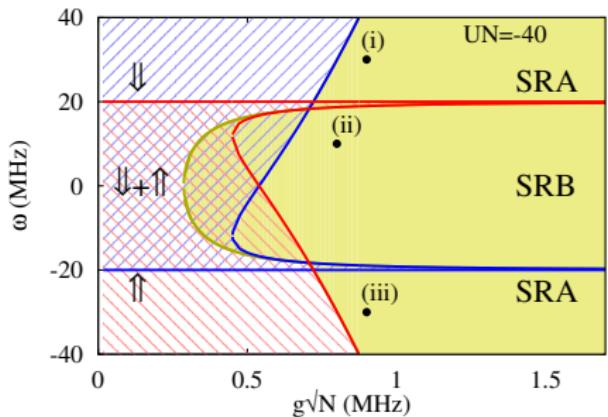
- Summary of experiment and classical dynamics
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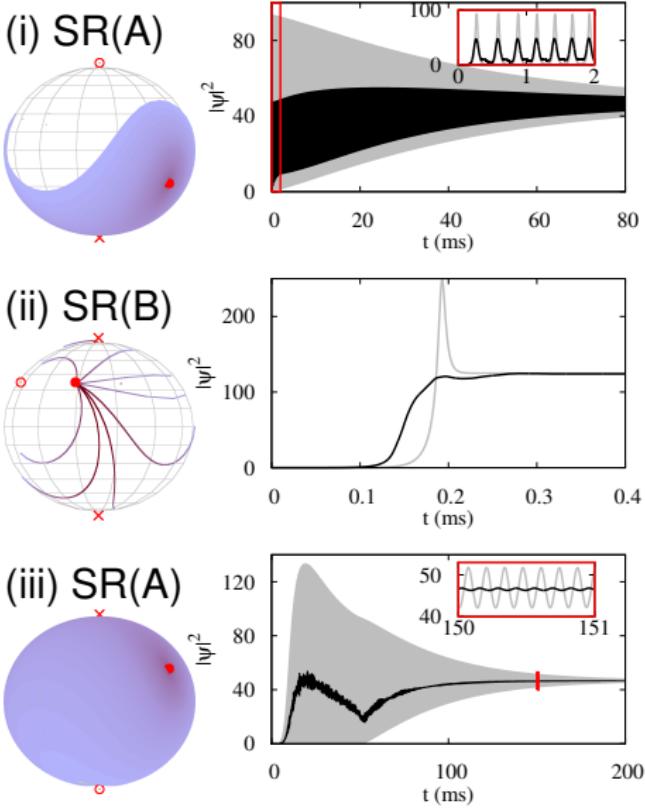
- JCHM vv Dicke
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- Disorder

Dynamics: Evolution from normal state

Gray: $\mathbf{S} = (\sqrt{N}, \sqrt{N}, -N/2)$
Black: Wigner distribution of \mathbf{S}, ψ



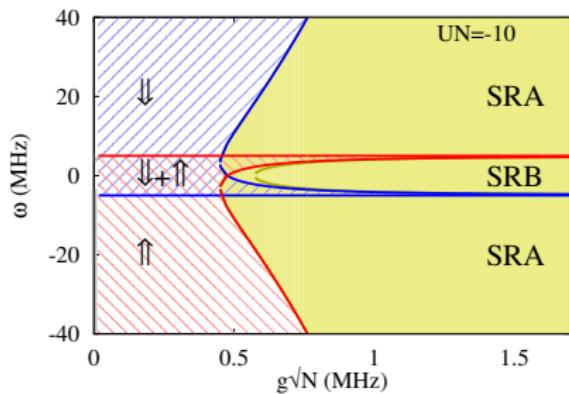
Oscillations: $\sim 0.1\text{ms}$
Decay: $20\text{ms}, 0.1\text{ms}, 20\text{ms}$



Asymptotic state: Evolution from normal state

(Near to experimental $UN = -13\text{MHz}$).

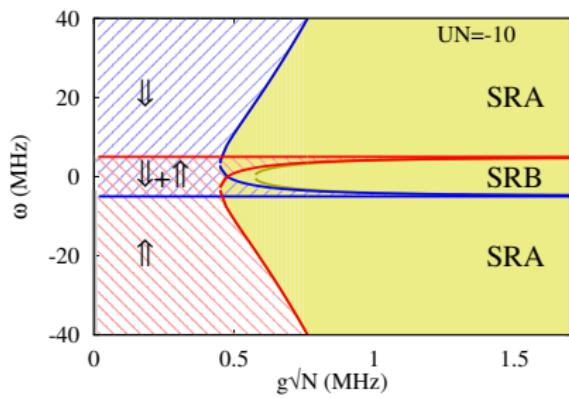
All stable attractors:



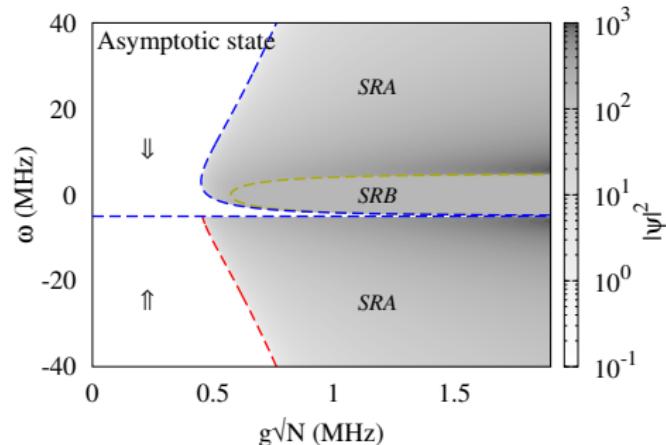
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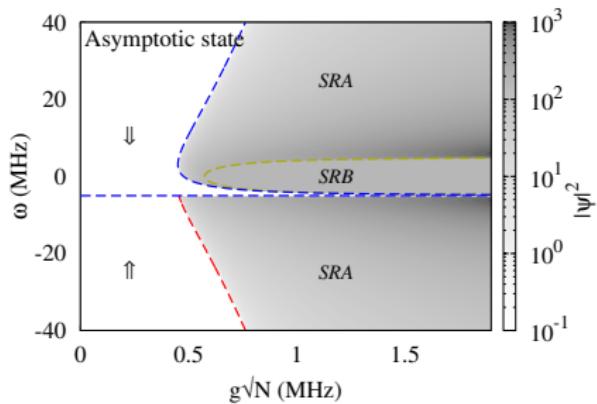
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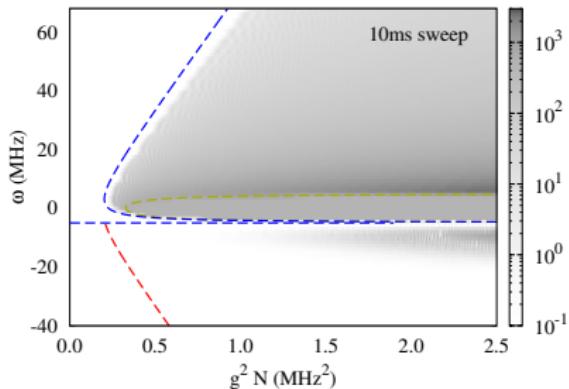
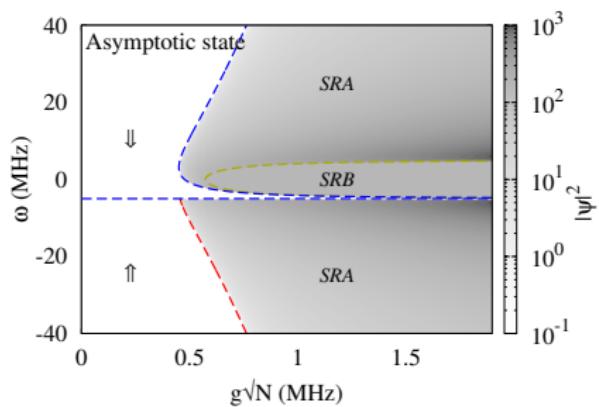
Starting from \downarrow



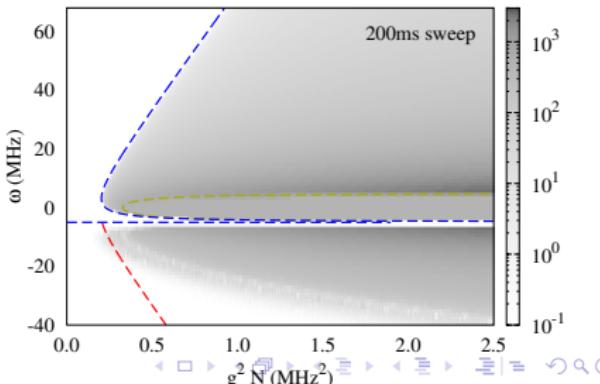
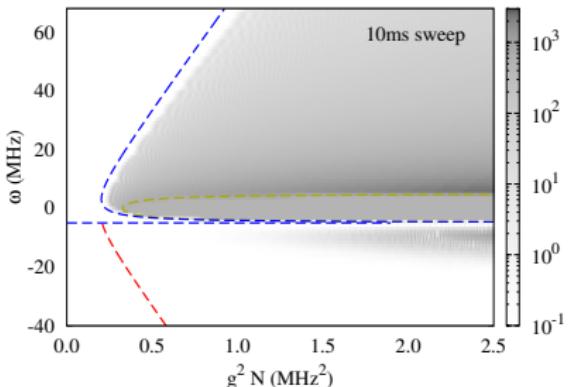
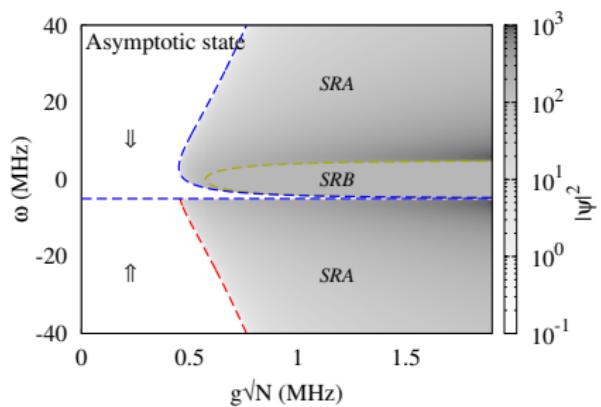
Timescales for dynamics: Consequences for experiment



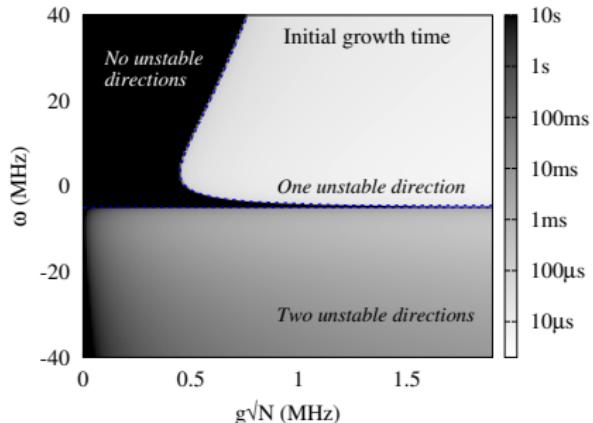
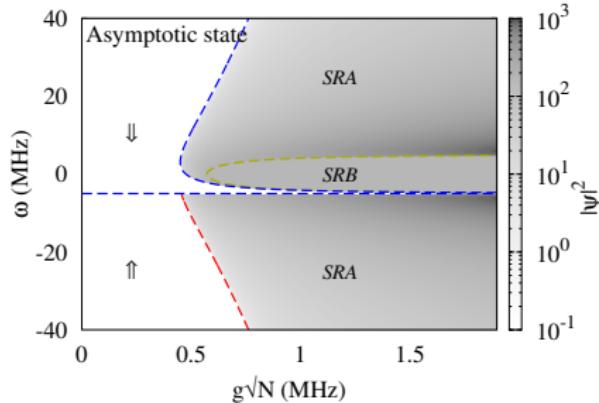
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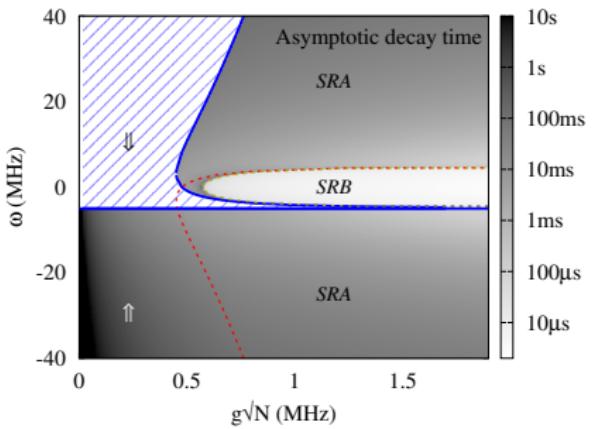
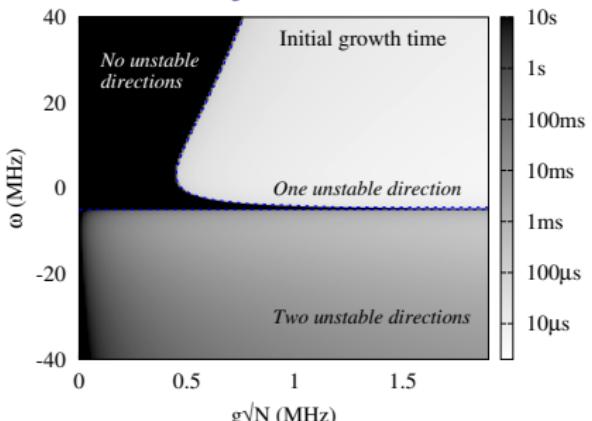
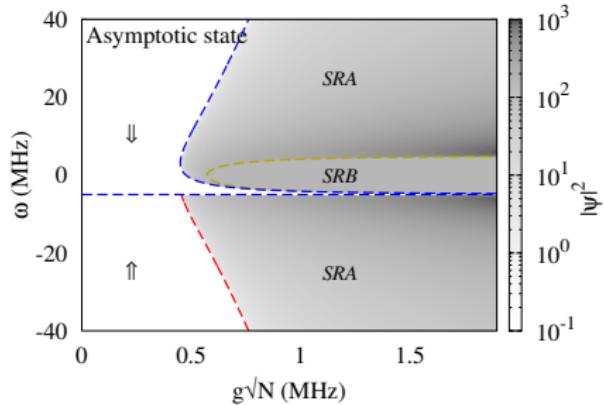
Timescales for dynamics: What are they?



Growth Most unstable eigenvalues near $\mathbf{S} = (0, 0, -N/2)$

Decay Slowest stable eigenvalues near final state

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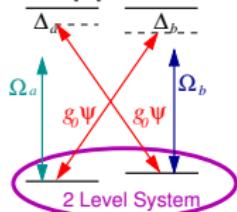


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Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ

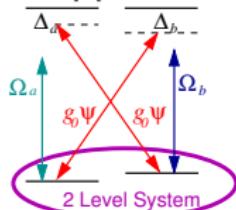


$$H = \dots + g(\psi^\dagger S^- + \psi S^+) + g'(\psi^\dagger S^+ + \psi S^-) + \dots$$

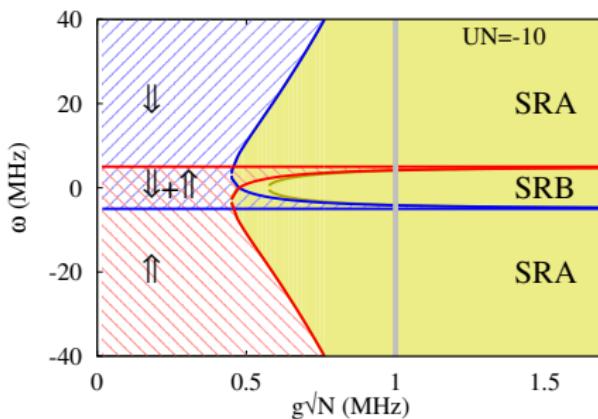
- SR(A) near phase boundary at small $\delta g \rightarrow$ Critical slowing down
- SR(A), SR(B) continuously connect

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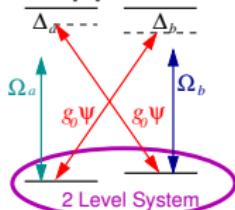
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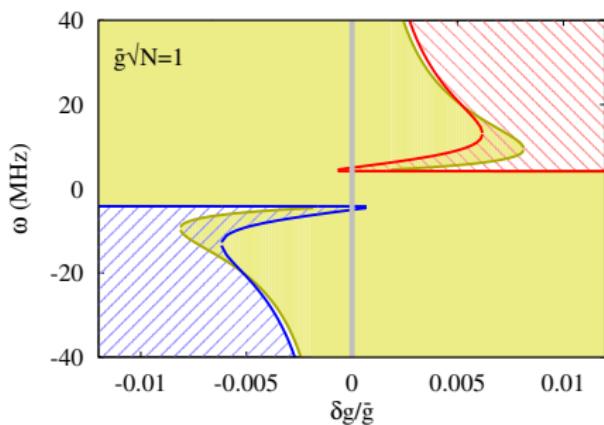
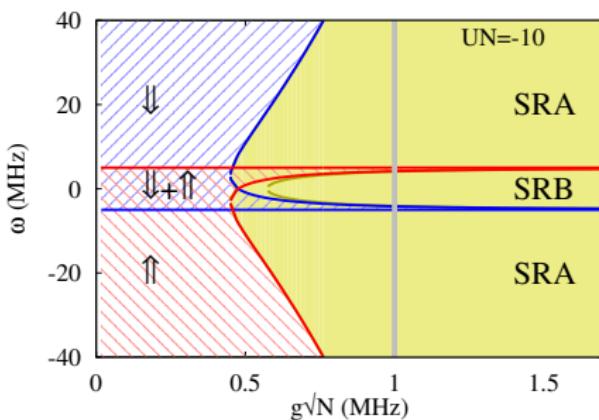
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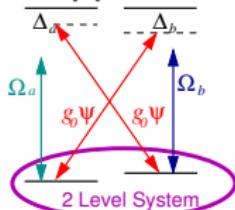
$$\delta g = g' - g, \quad 2\bar{g} = g' + g$$



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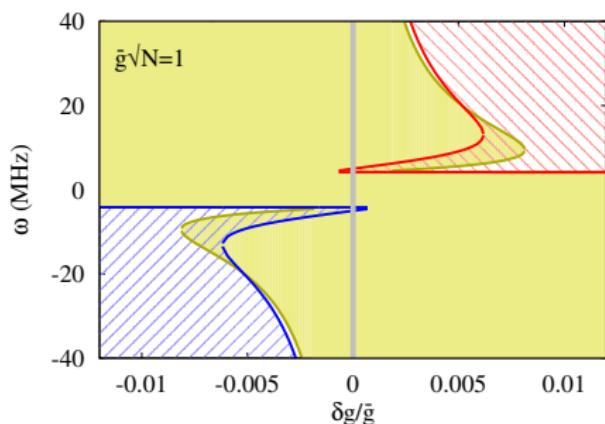
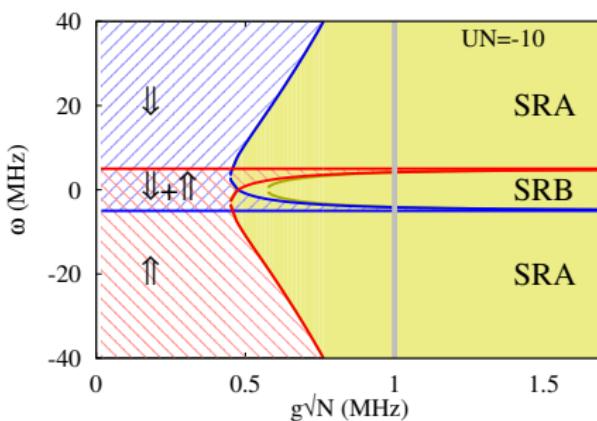
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Dynamics of generalized Dicke model



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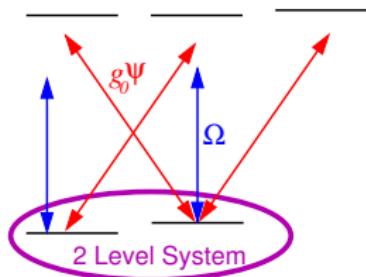
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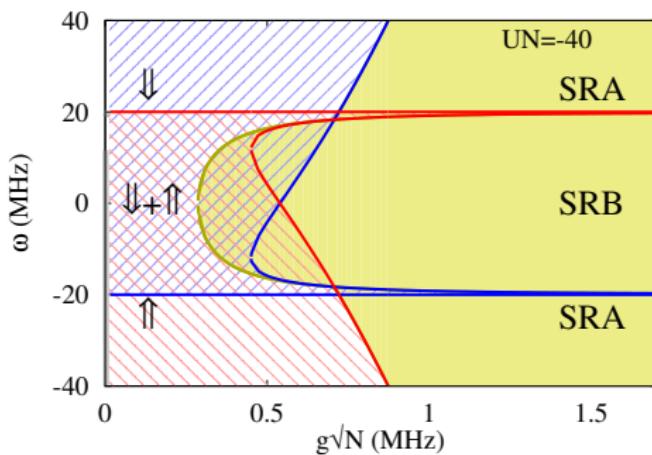
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Regions without fixed points

Changing U :

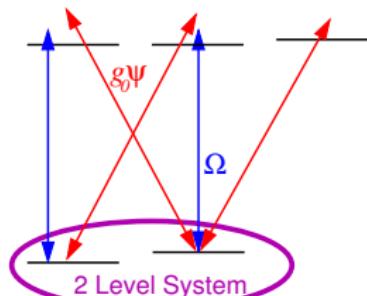


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

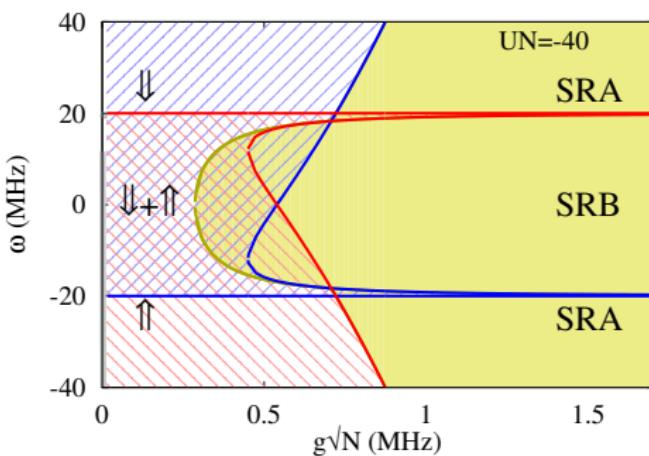


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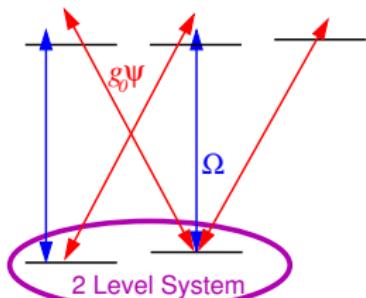


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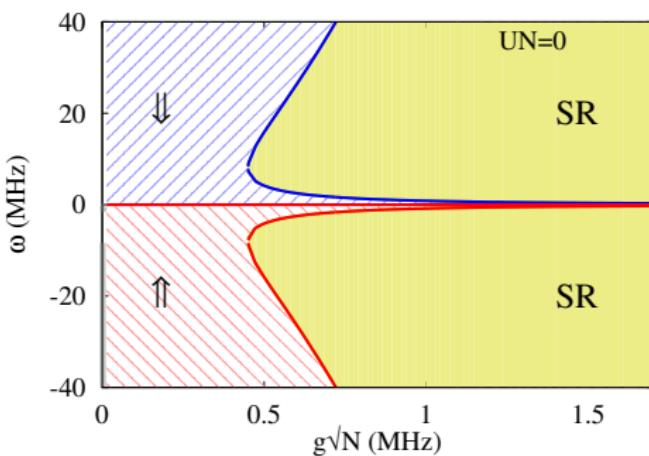


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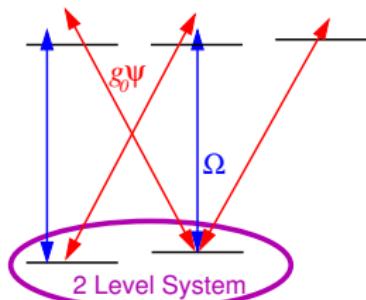


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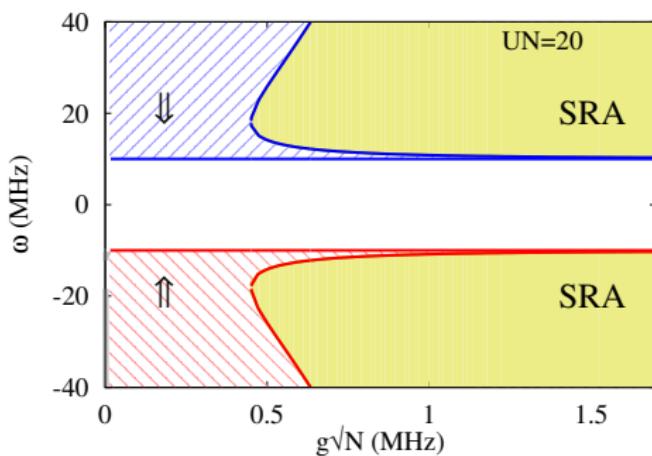


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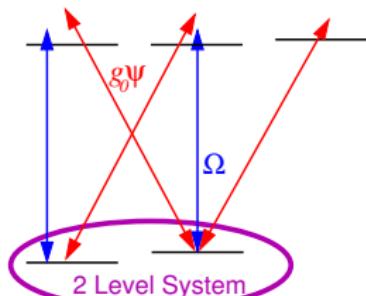


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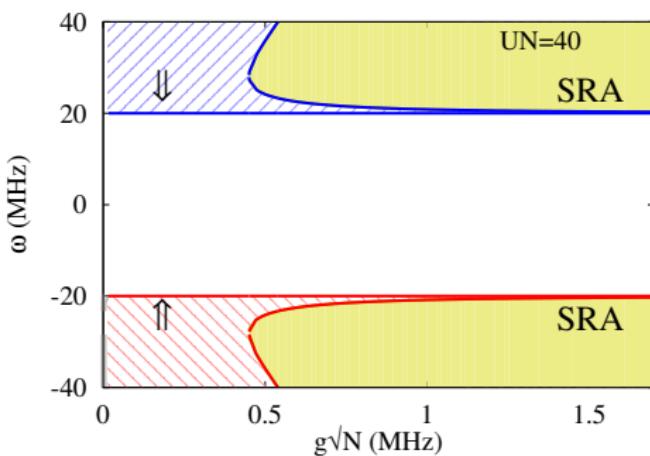


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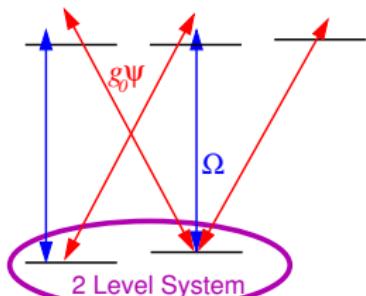


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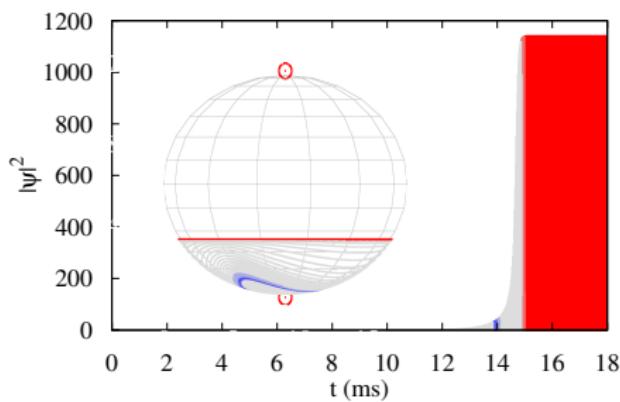
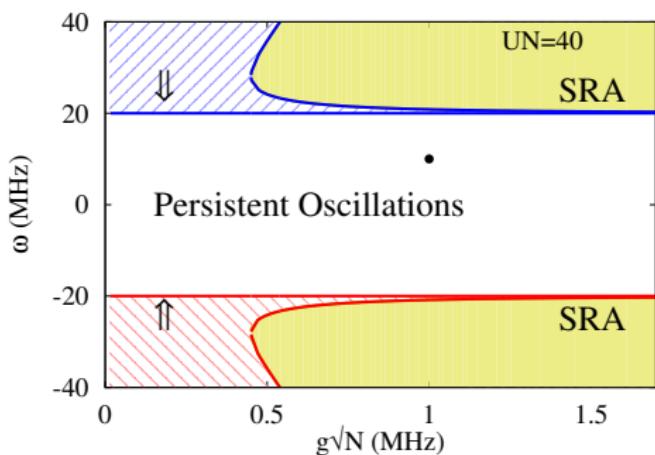


Regions without fixed points

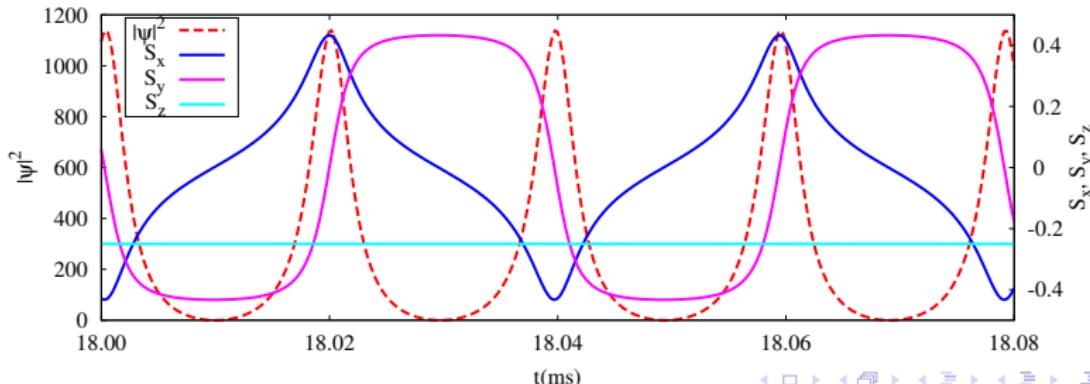
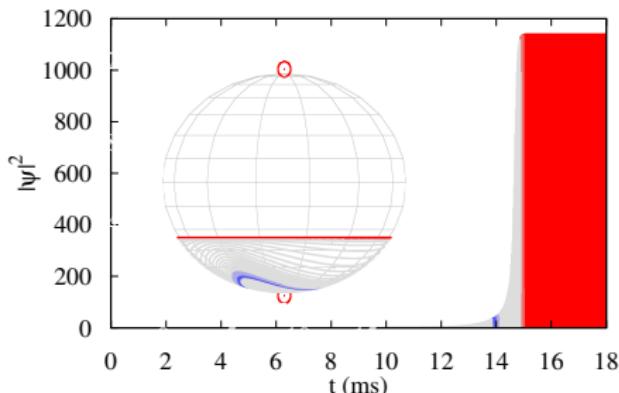
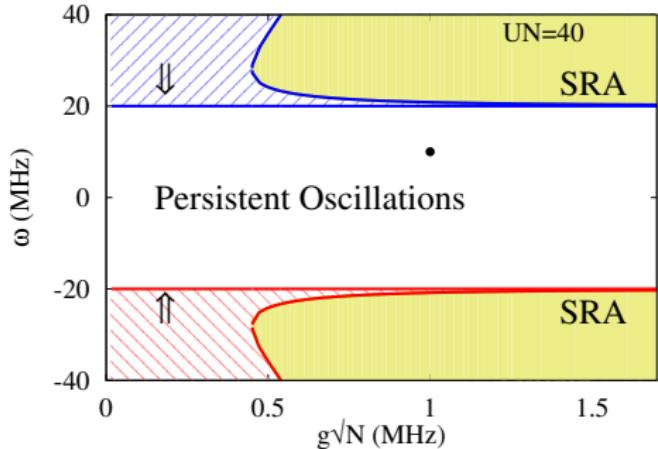
Changing U :



$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$



Persistent (optomechanical) oscillations



Jaynes Cummings Hubbard model



1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

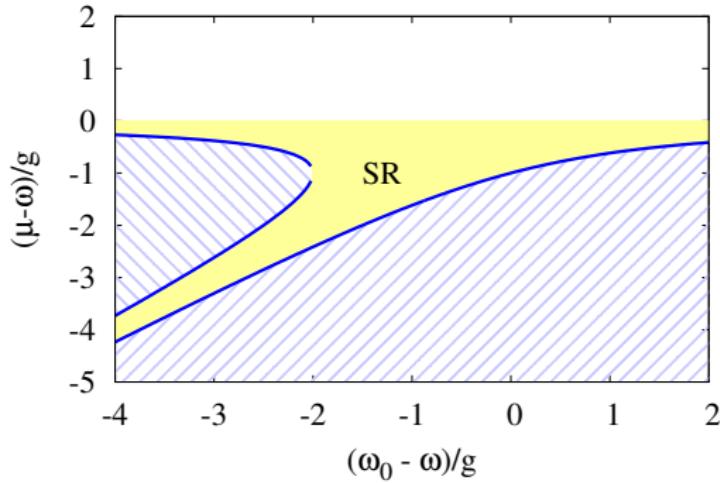
- Summary of experiment and classical dynamics
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- Persistent oscillating phases

3 Jaynes Cummings Hubbard model

- JCHM vv Dicke
- Coherently driven array
- Disorder

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$

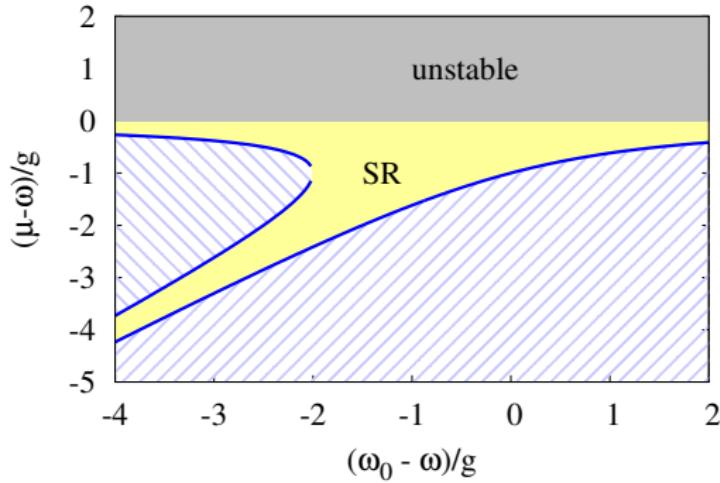


- Transition at:
 $g^2 N > (\omega - \mu)|\omega_0 - \mu|$
- Reduce critical g

[Eastham and Littlewood, PRB '01]

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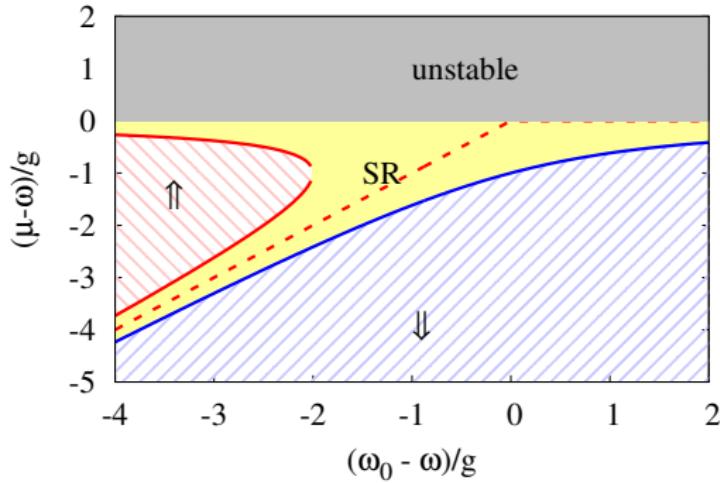
- Transition at: $g^2N > (\omega - \mu)|\omega_0 - \mu|$
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- Unstable if $\mu > \omega$

Inverted if $\mu > \omega_0$

[Eastham and Littlewood, PRB '01]

Equilibrium: Dicke model with chemical potential

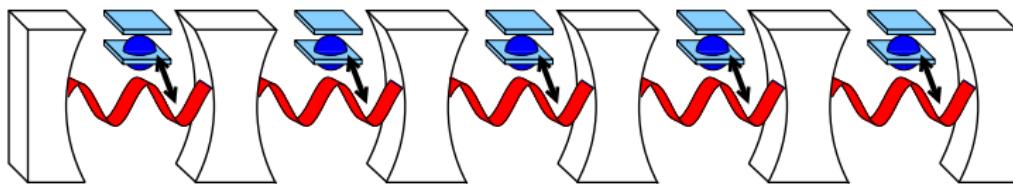
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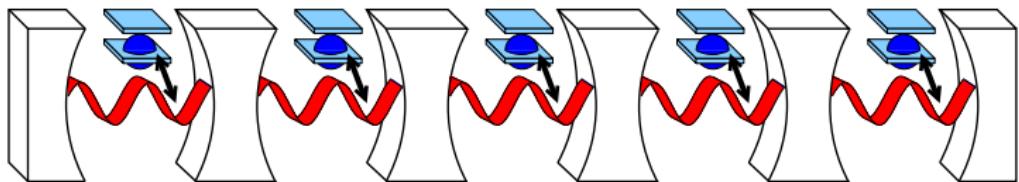
[Eastham and Littlewood, PRB '01]

Jaynes-Cummings Hubbard model

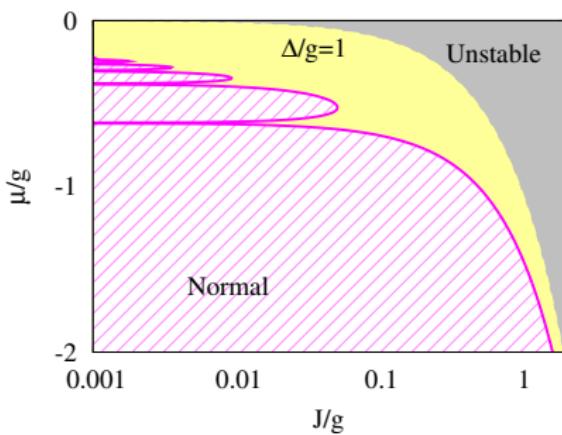


$$H = -\frac{J}{z} \sum_{ij} \psi_i^\dagger \psi_j + \sum_i \frac{\Delta}{2} \sigma_i^z + g(\psi_i^\dagger \sigma_i^- + \text{H.c.})$$

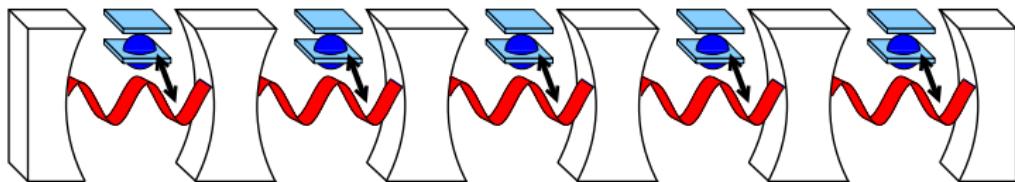
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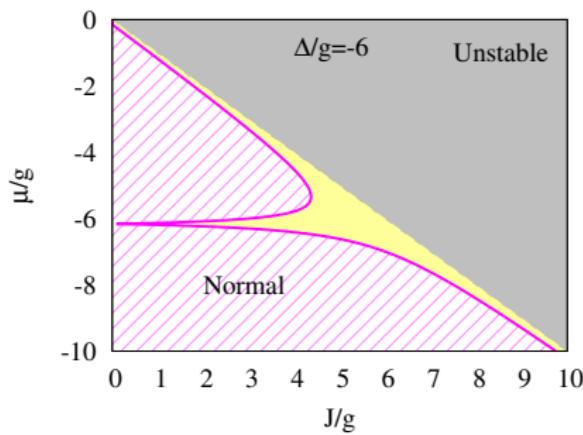
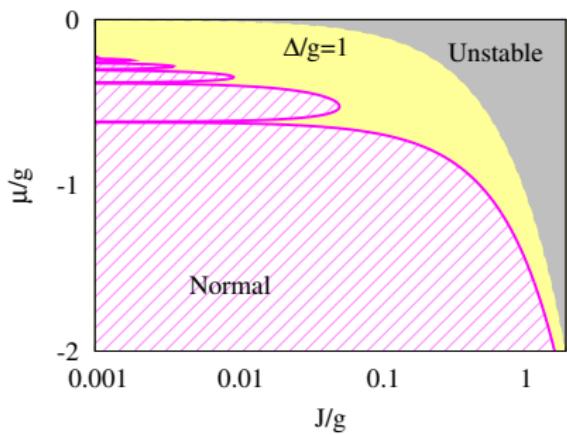
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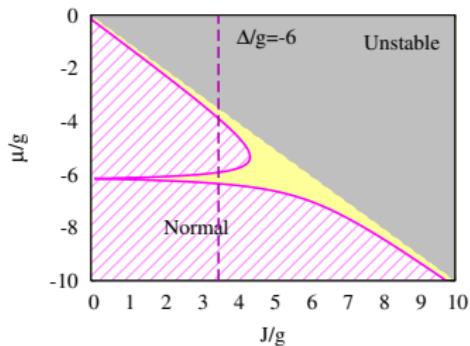


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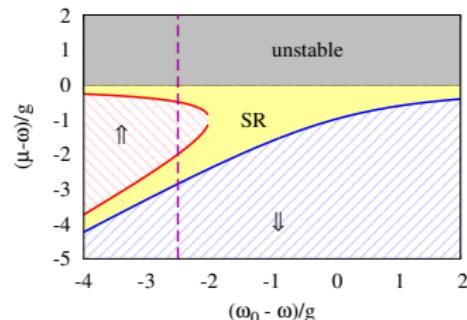


Dicke vs JCHM

JCHM



Dicke

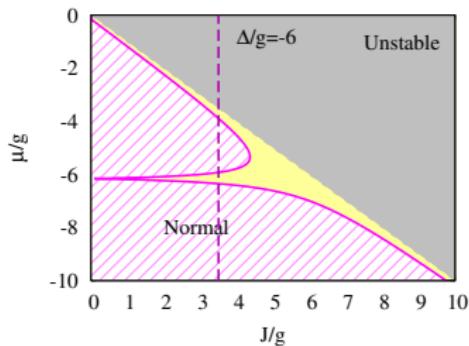


\rightarrow $\downarrow\downarrow$ \rightarrow $\uparrow\downarrow$ \rightarrow Dicke photon mode

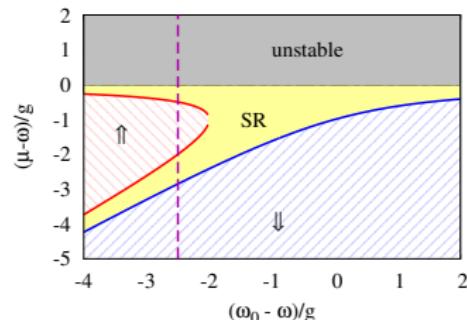
\rightarrow $\uparrow\downarrow$ \rightarrow $\square = 1$ Mott lobe

Dicke vs JCHM

JCHM



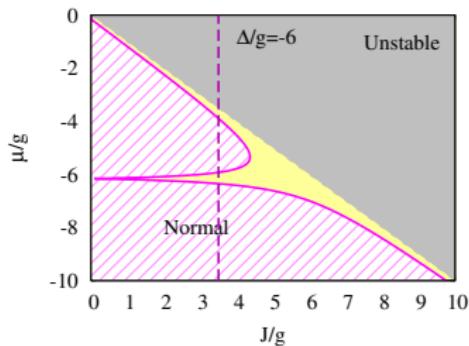
Dicke



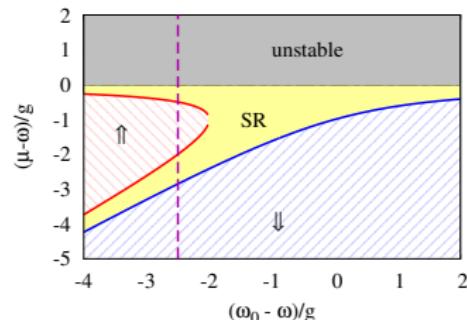
- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode

Dicke vs JCHM

JCHM

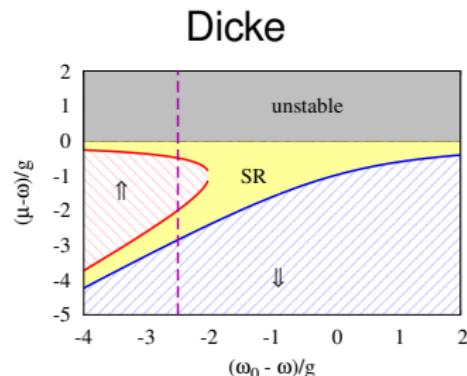
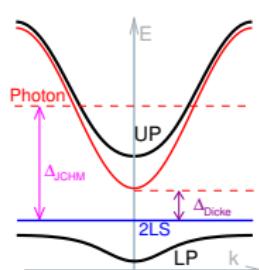
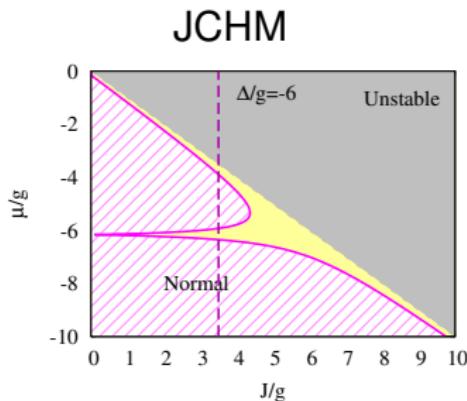


Dicke



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- $\uparrow \leftrightarrow n = 1$ Mott lobe

Dicke vs JCHM



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Jaynes Cummings Hubbard model



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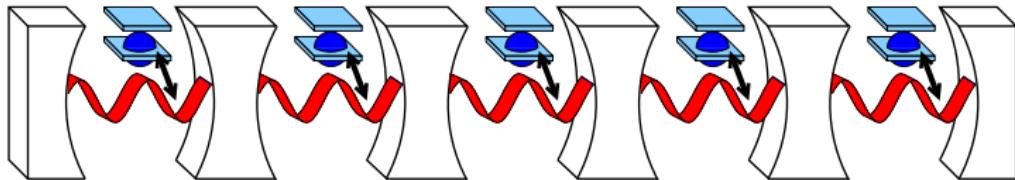
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- **Coherently driven array**
- Disorder

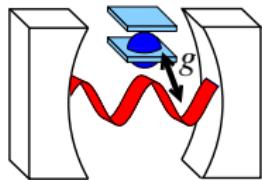
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$$\partial_t \rho = -i[H, \rho] - \frac{\kappa}{2} L_\psi[\rho] - \frac{\gamma}{2} L_{\sigma^-}[\rho]$$

Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]

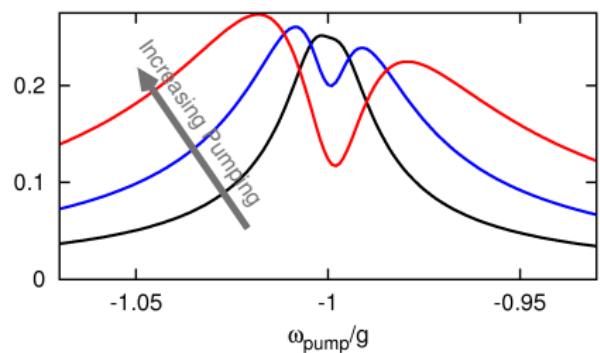


$$H = \frac{\Delta}{2} \sigma^z + g(\psi^\dagger \sigma^- + \text{H.c.}) + f(\psi e^{i\omega_{\text{pump}} t} + \text{H.c.})$$
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Anti-resonance in $\langle \hat{a} \rangle$

Fluorescence

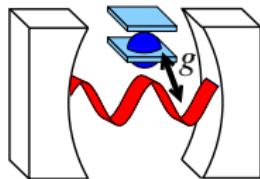
Fluorescence intensity



Mollow triplet fluorescence

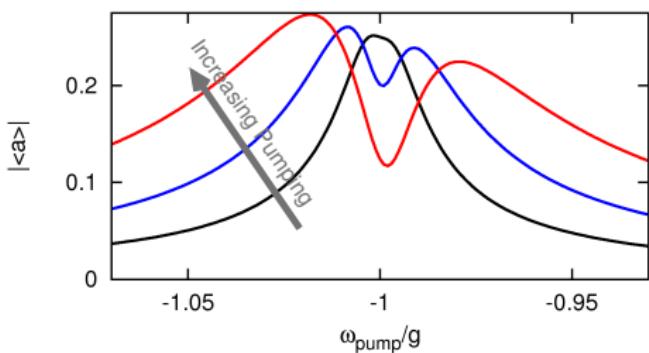
[Lang *et al.* PRL '11]

Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



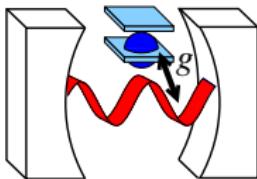
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- Anti-resonance in $|\langle \psi \rangle|$.
- Effective 2LS:
 $|\text{Empty}\rangle, |\text{1 polariton}\rangle$



[Lang *et al.* PRL '11]

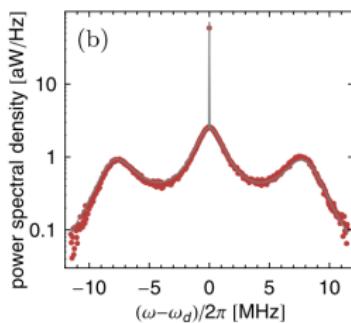
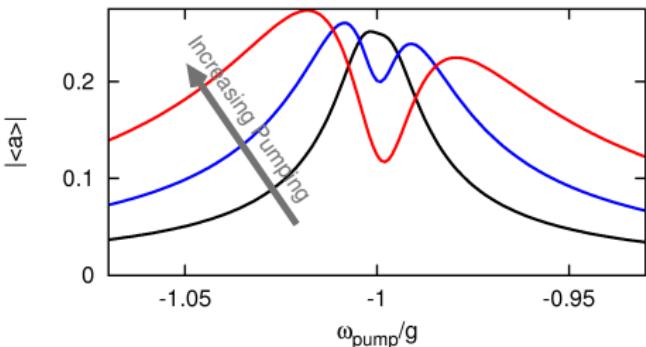
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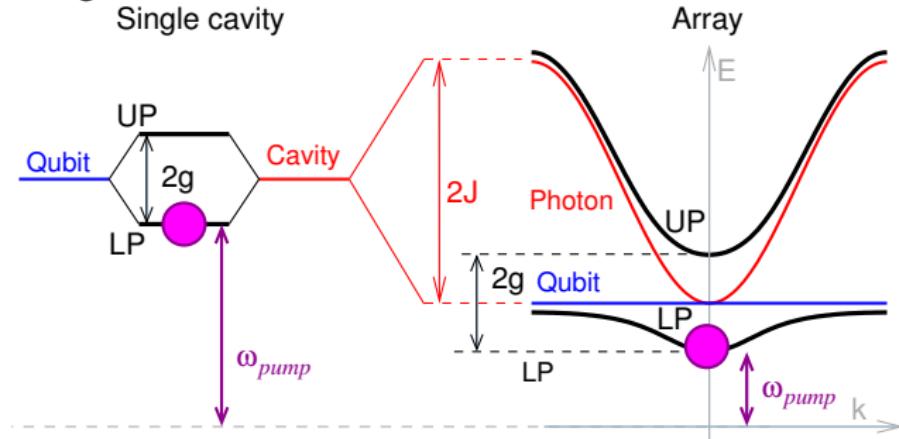
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[Lang *et al.* PRL '11]

Coherently pumped dimer & array

Chose detuning *a la* Dicke model

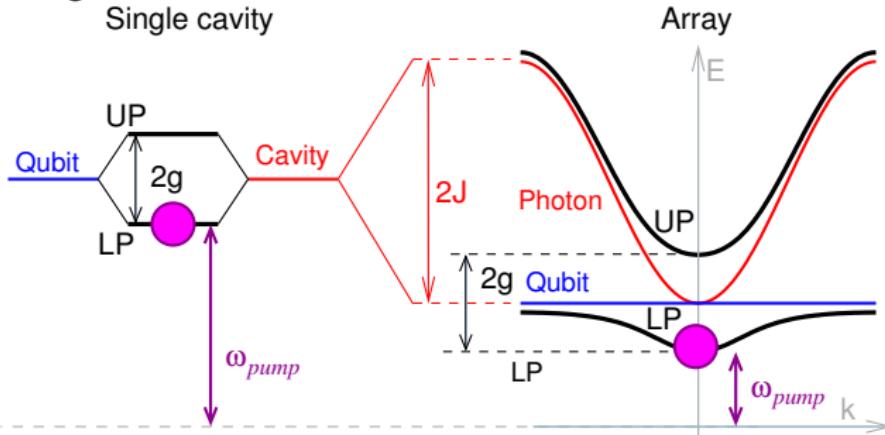


- Bistability at intermediate J
- More/less localised states
- Connection to Dicke limit

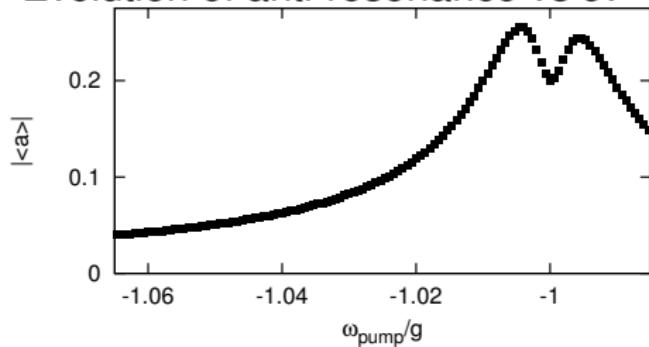
[Nissen *et al.* PRL '12]

Coherently pumped dimer & array

Chose detuning *a la* Dicke model



Evolution of anti-resonance vs J .

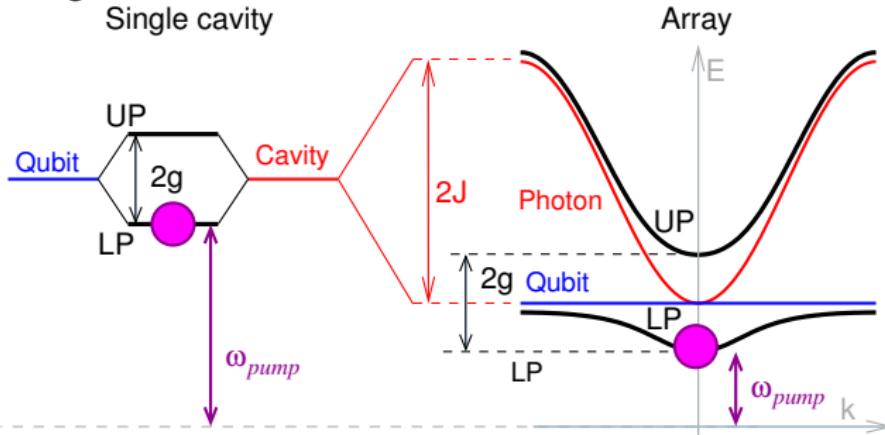


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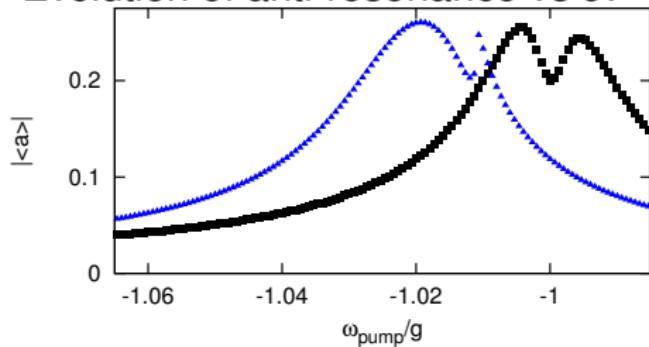
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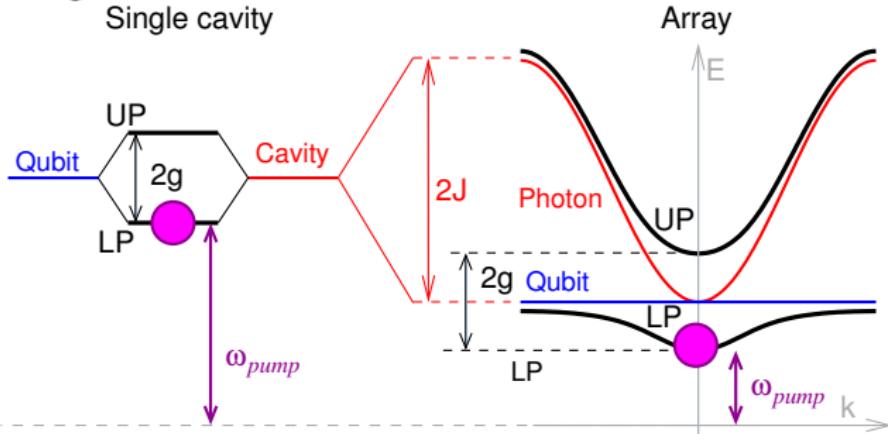


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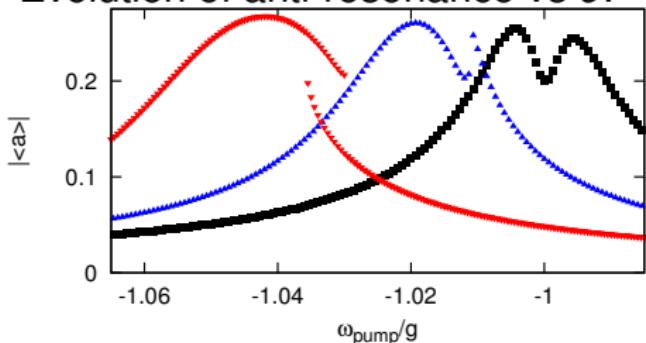
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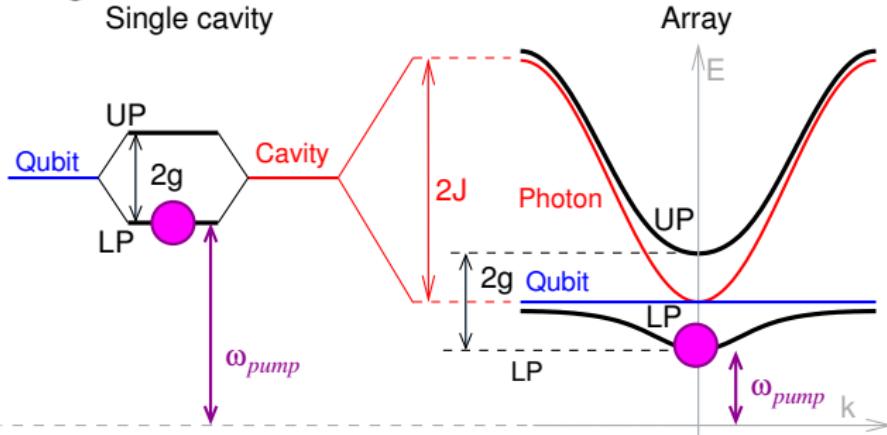


Bistability at intermediate J
More/less localised states
Approaching Dicke limit

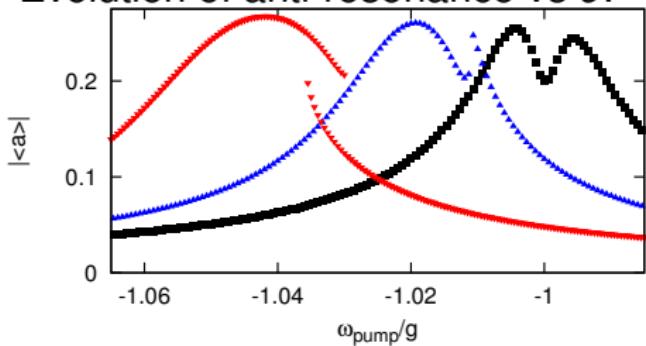
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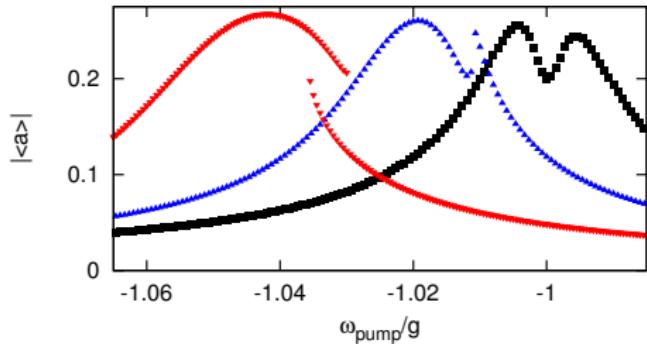
- Bistability at intermediate J
 - ▶ More/less localised states
 - ▶ Connects to Dicke limit

[Nissen *et al.* PRL '12]

Photon blockade picture $J \lesssim g$

- Polariton basis
- Nonlinearity $|\epsilon_2 - 2\epsilon_1| \propto g$.

$$H = \sum_i \left(\frac{\epsilon}{2} \tau_i^z + \tilde{f} \tau_i^x \right)$$



[Nissen *et al.* PRL '12]

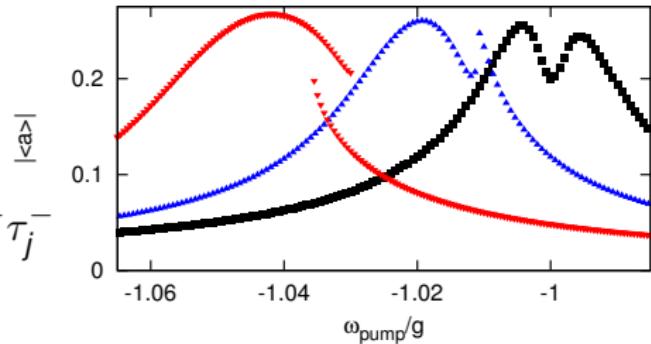
- Decouple hopping:
 $\tau^z \tau^z \rightarrow \tau^+ \tau^- + \tau^- \tau^+$
- Bistability for

$$J > J_c = \frac{4}{P} \left(\frac{2B + (g/2)^2}{3} \right)^{3/2}$$

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[Nissen et al. PRL '12]

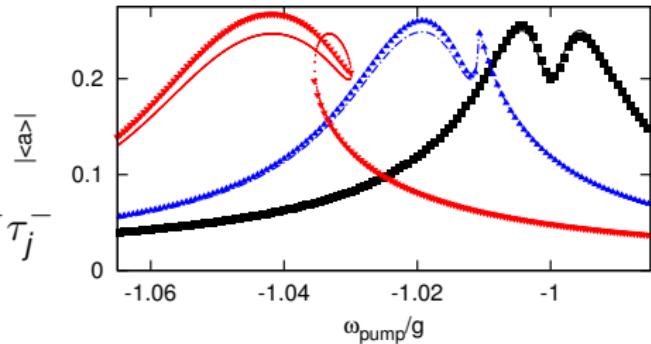
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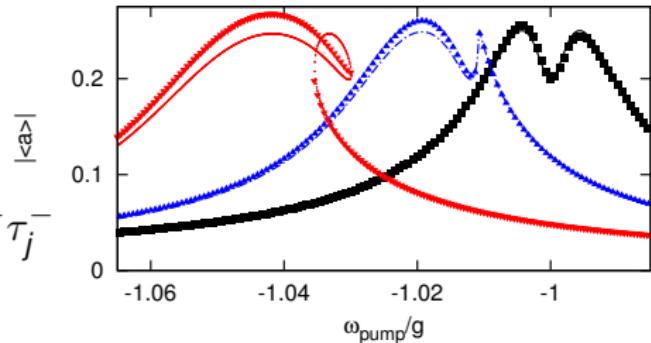
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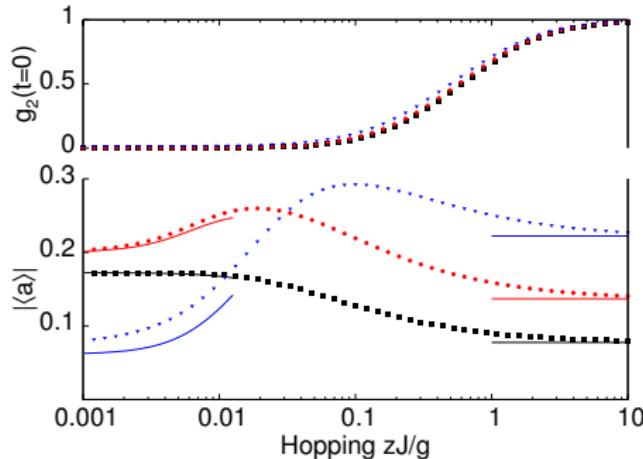


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Coherently pumped array: correlations & fluorescence

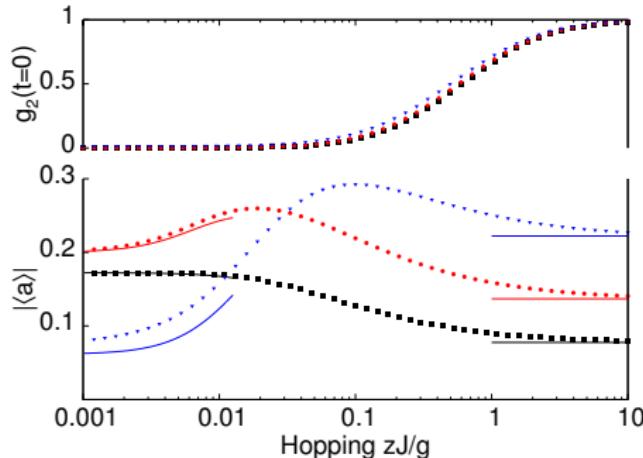


Correlations

- $g_2 : 0 \rightarrow 1$ crossover.

- Small J: Mollow triplet
- Large J: Off resonance fluorescence
- Pump at collective resonance
- Mismatch if $J \neq 0$

Coherently pumped array: correlations & fluorescence

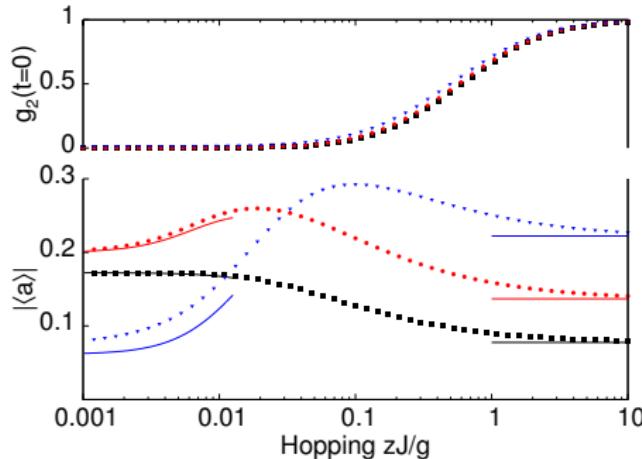


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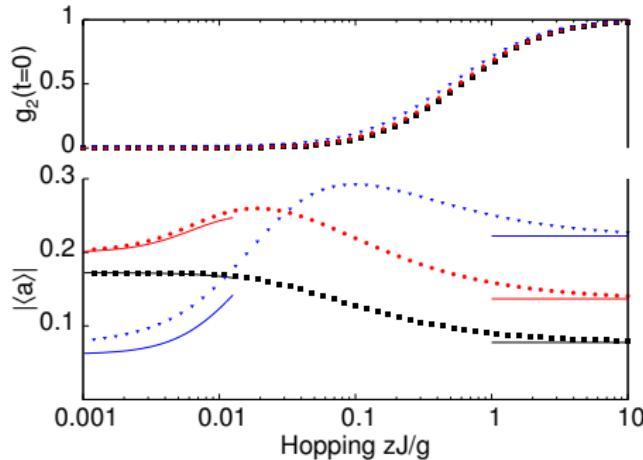
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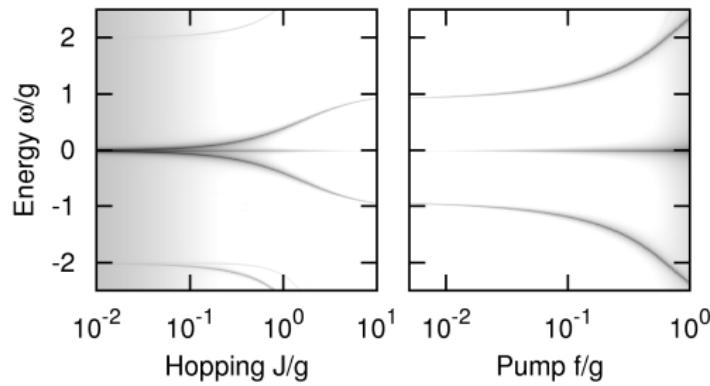


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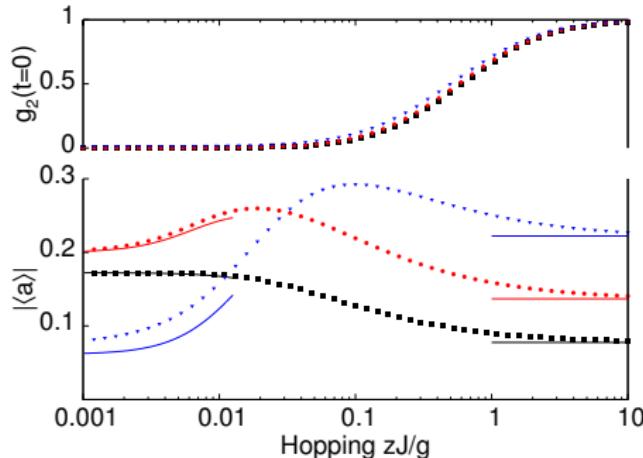
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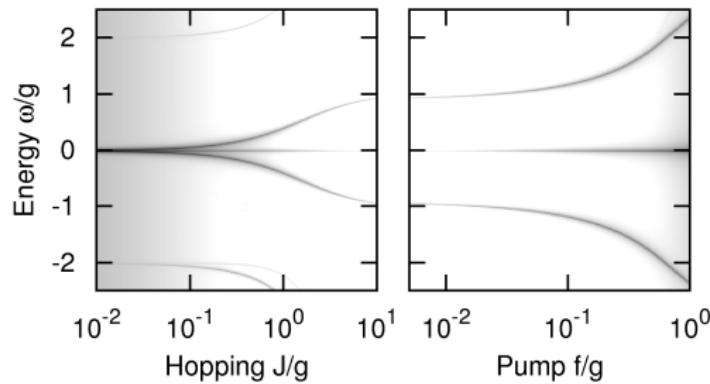


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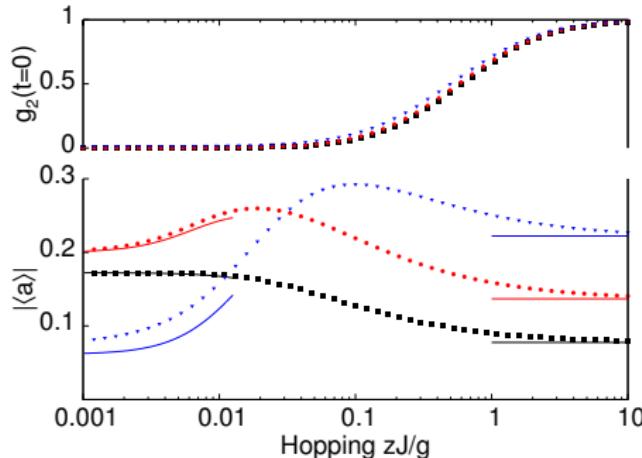
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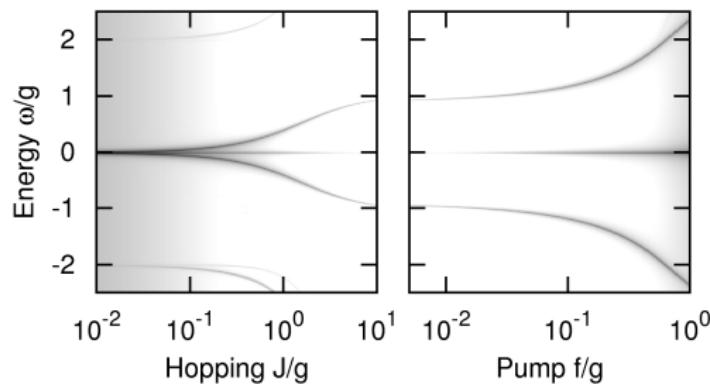


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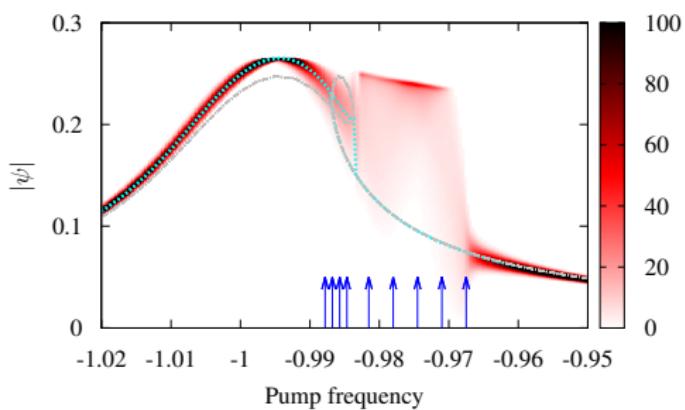
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Coherent pumped array – disorder

- Effect of disorder, $\Delta \rightarrow \Delta_i$
 - ▶ Distribution of ψ – Washes out bistable jump

- Bistability disappears → phase transition on Δ_i
- Complex ψ distribution
- Superfluid phases in driven system?



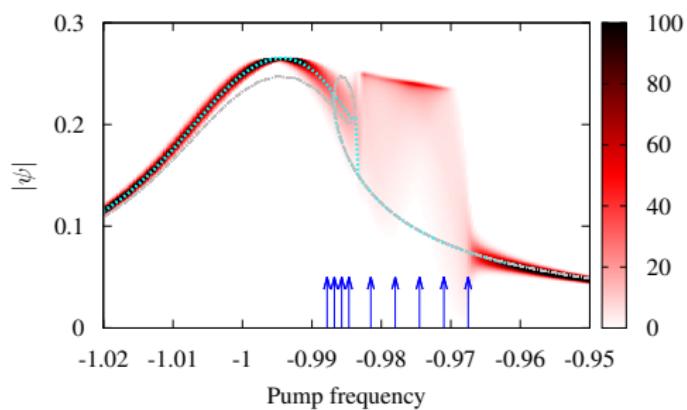
[Kulaitis *et al.* PRA, '13]

Coherent pumped array – disorder

- Effect of disorder, $\Delta \rightarrow \Delta_i$
 - ▶ Distribution of ψ – Washes out bistable jump
- Bistability near resonance — phase of ψ depends on Δ_i

• Complex ψ distribution

• Superfluid phases in driven system?

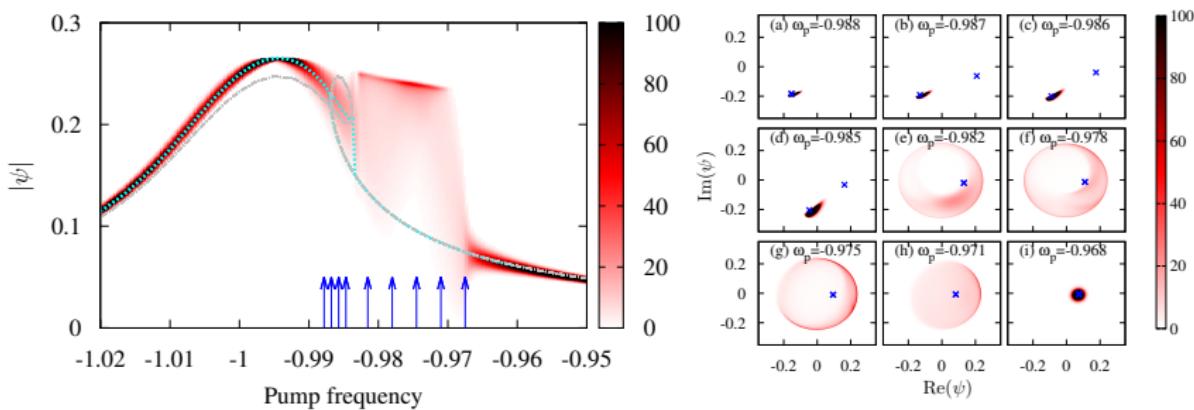


[Kulaitis *et al.* PRA, '13]

Coherent pumped array – disorder

- Effect of disorder, $\Delta \rightarrow \Delta_i$
 - ▶ Distribution of ψ – Washes out bistable jump
- Bistability near resonance — phase of ψ depends on Δ_i
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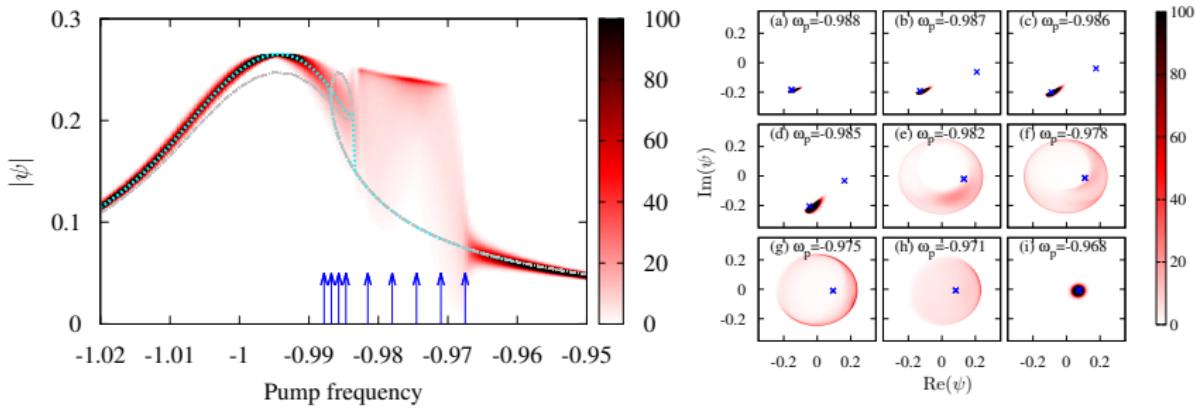
Is there more to say about driven system?



[Kulaitis et al. PRA, '13]

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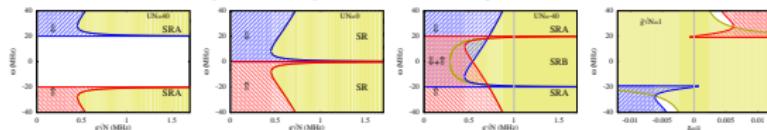
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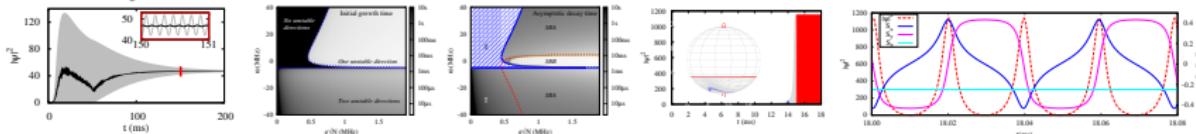
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Summary

- Wide variety of dynamical phases

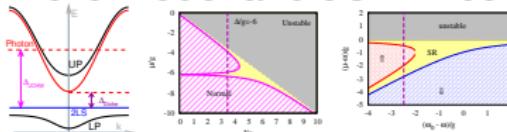


- Slow dynamics for $U < 0$ & Persistent oscillations for $U > 0$

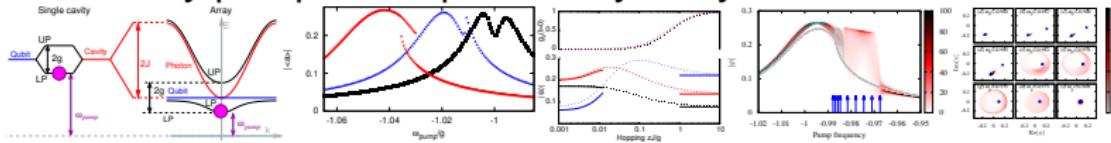


JK et al. PRL '10, Bhaseen et al. PRA '12

- Dicke model and JCHM: connection at $J \rightarrow \infty$



- Coherently pumped coupled cavity array



Nissen et al. PRL '12, Kulaitis et al. PRA '13

4 Ferroelectric transition

5 Dicke vs JCHM

6 Pumping without symmetry breaking

7 Collective dephasing

Ferroelectric transition

Atoms in Coulomb gauge

$$H = \sum \omega_k a_k^\dagger a_k + \sum_i [p_i - eA(r_i)]^2 + V_{\text{coul}}$$

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$$H = \omega_0 S^z + \omega \psi^\dagger \psi + g(S^+ + S^-)(\psi + \psi^\dagger) + N\zeta(\psi + \psi^\dagger)^2 - \eta(S^+ - S^-)^2$$

(nb $g^2, \zeta, \eta \propto 1/V$).

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Gauge transform to dipole gauge $\mathbf{D} \cdot \mathbf{r}$

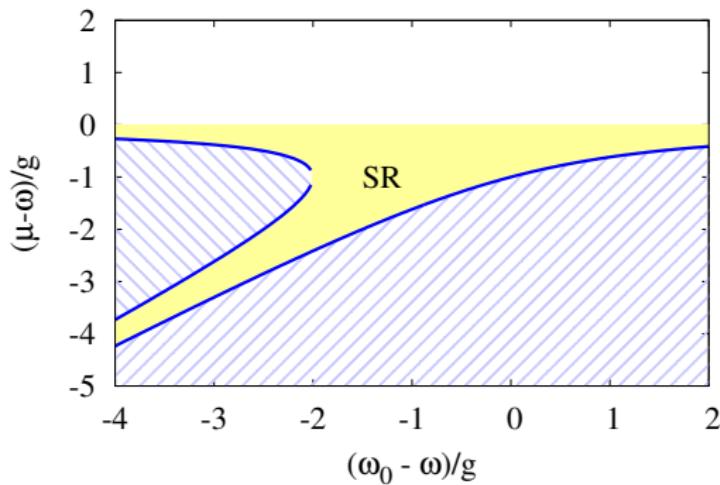
$$H = \omega_0 S^z + \omega \psi^\dagger \psi + \bar{g}(S^+ - S^-)(\psi - \psi^\dagger)$$

“Dicke” transition at $\omega_0 < N\bar{g}^2/\omega \equiv 2\eta N$

But, ψ describes electric displacement

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$

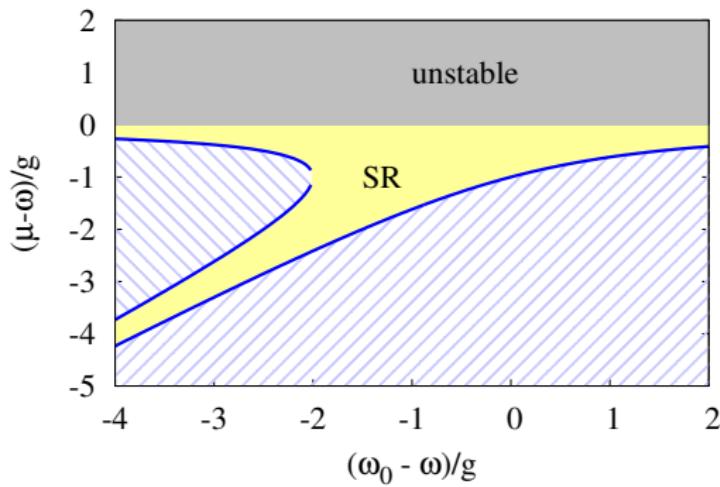


- Transition at:
 $g^2 N > (\omega - \mu)|\omega_0 - \mu|$
- Reduce critical g

[Eastham and Littlewood, PRB '01]

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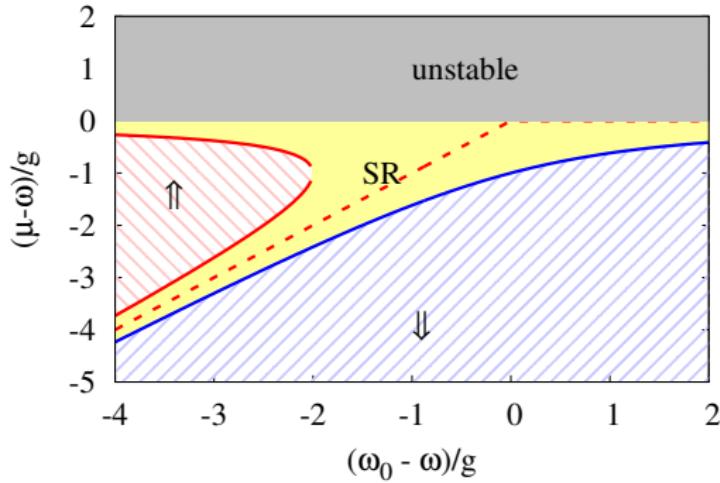
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Inverted if $\mu > \omega_0$

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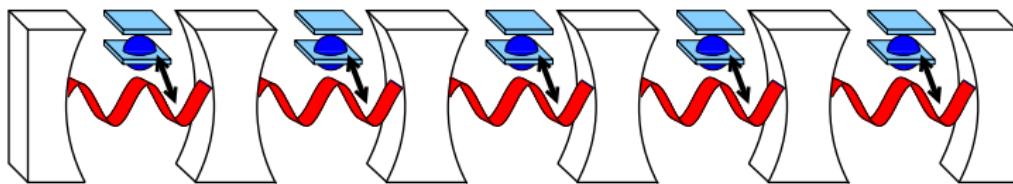
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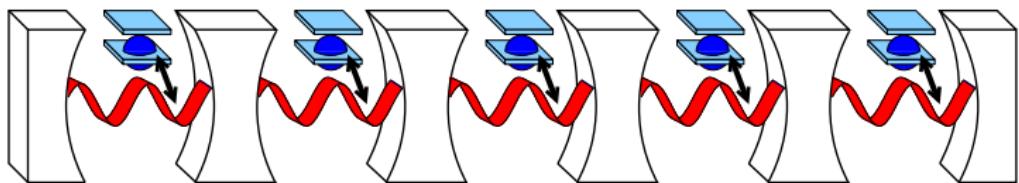
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Jaynes-Cummings Hubbard model

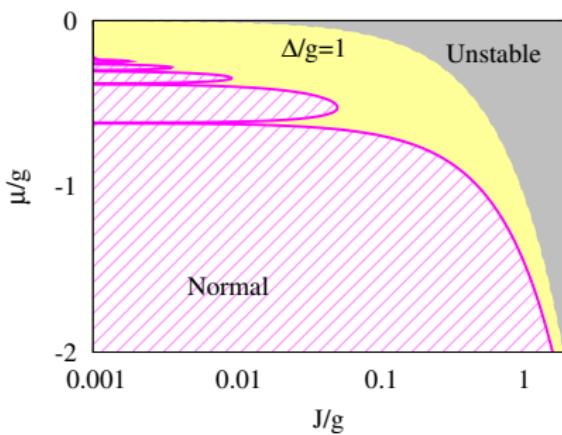


$$H = -\frac{J}{z} \sum_{ij} \psi_i^\dagger \psi_j + \sum_i \frac{\Delta}{2} \sigma_i^z + g(\psi_i^\dagger \sigma_i^- + \text{H.c.})$$

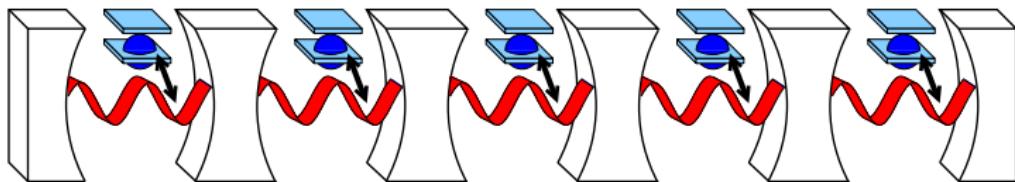
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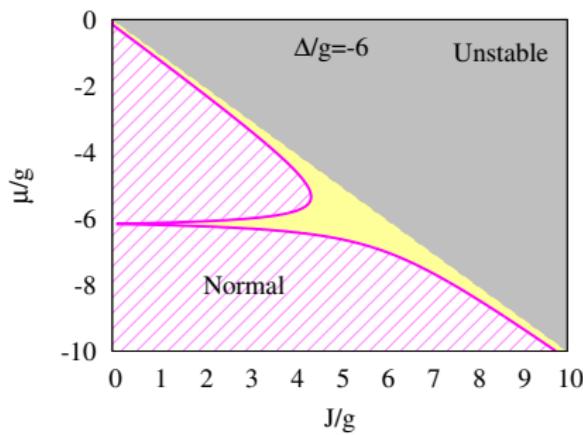
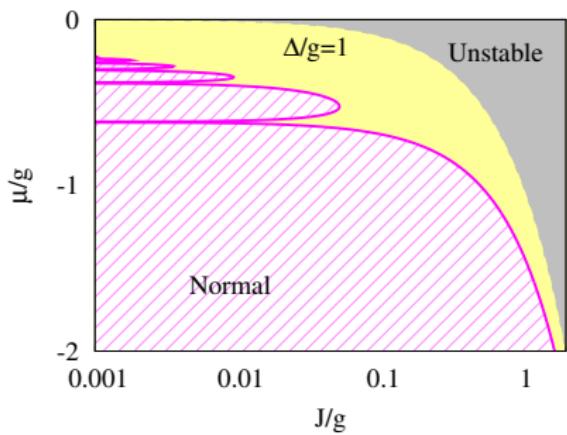
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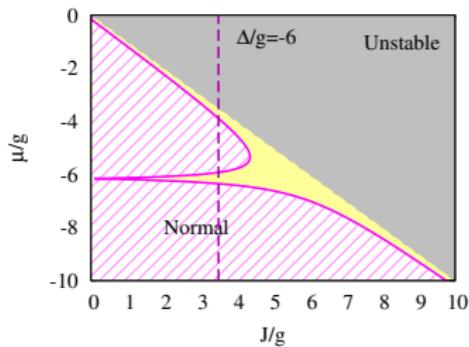


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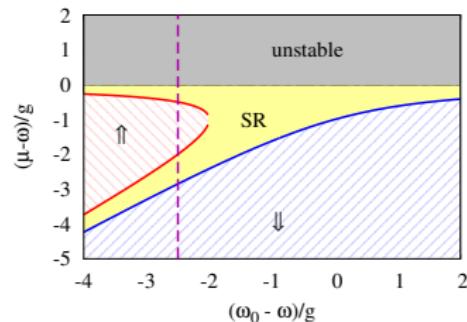


Dicke vs JCHM

JCHM



Dicke

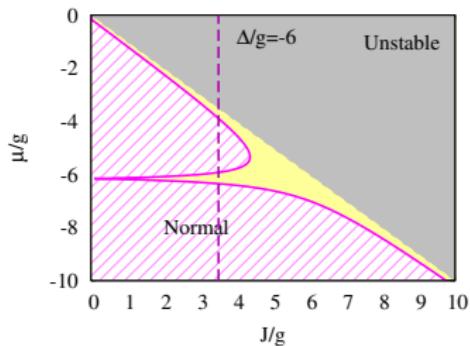


\rightarrow $\omega_0 - \omega = 0$ case of JCHM \rightarrow Dicke photon mode

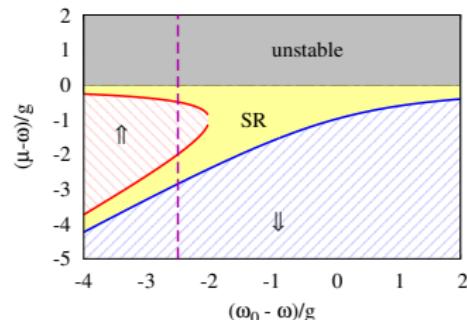
\rightarrow \uparrow \leftrightarrow \downarrow = 1 Mott lobe

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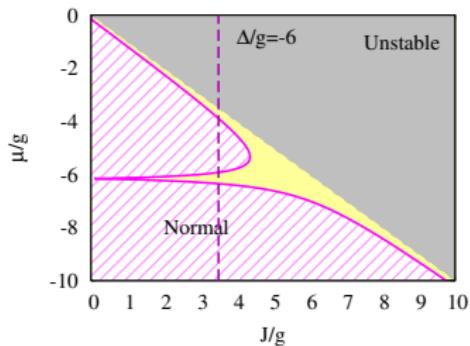
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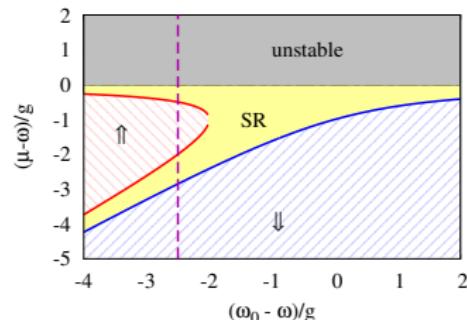
- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode

Dicke vs JCHM

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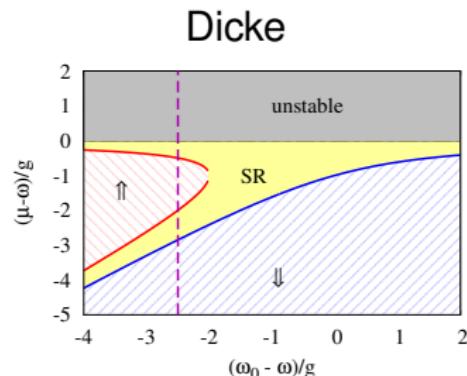
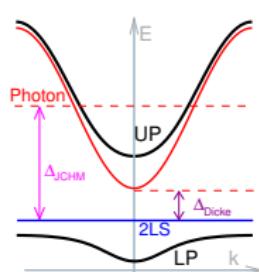
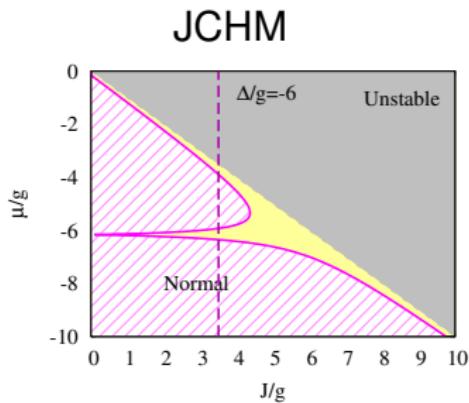


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Raman pumping

- How to pump without breaking symmetry
- Counter-rotating terms — Raman pumping
 - ▶ Atom proposal [Dimer *et al.* PRA '07]
 - ▶ Atom experiment [Baumann *et al.* Nature '10]

JK, Türeci, Houck in progress

• Qubit dephasing much bigger than atom

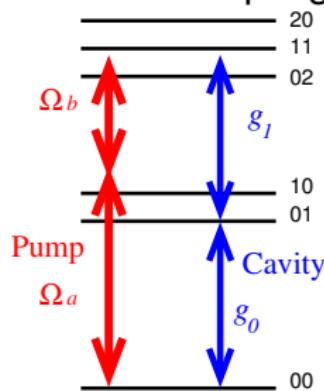
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Tunable-coupling-qubit

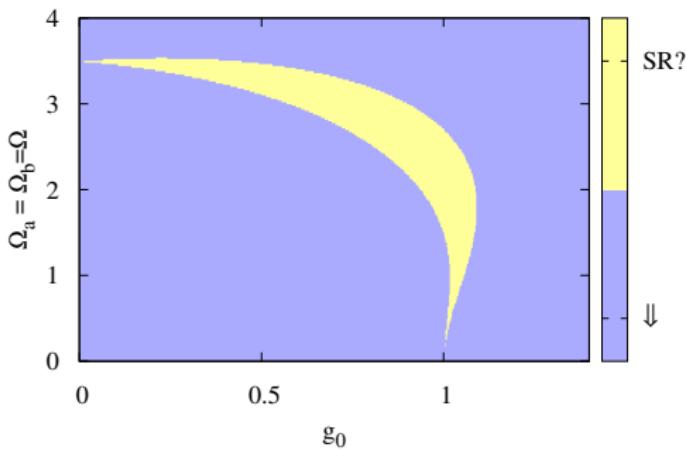
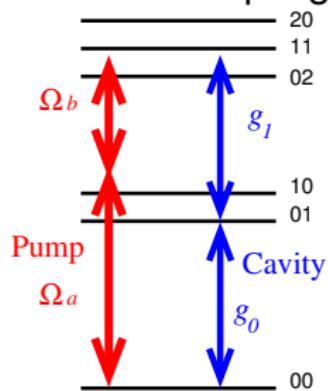


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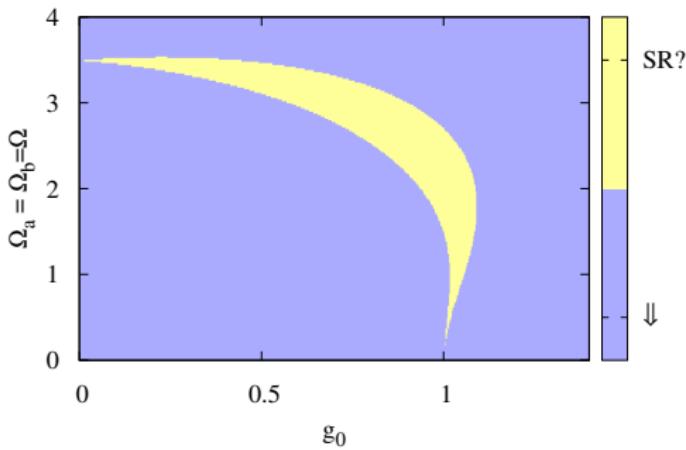
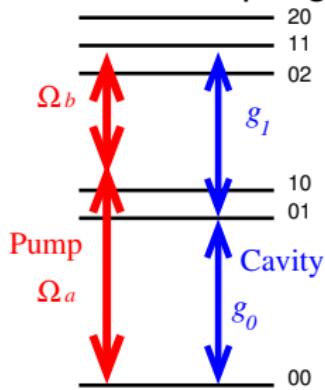


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Collective dephasing

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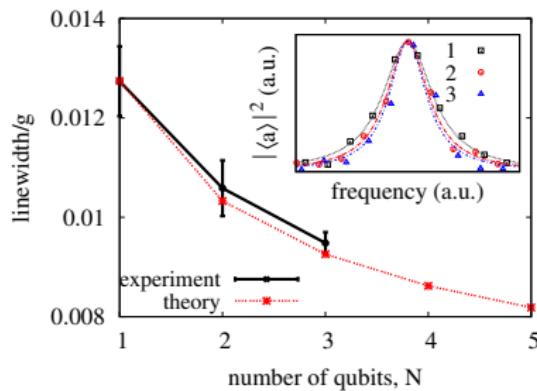
[Nissen, Fink *et al.* arXiv:1302.0665]

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